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## Existence of a Renormalized Solution of Nonlinear Parabolic Equations with Lower Order Term and General Measure Data

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ABSTRACT: We give an existence result of a renormalized solution for a class of nonlinear parabolic equations  $\frac{\partial b(u)}{\partial t} - \operatorname{div} \left( a(x,t,\nabla u) \right) + H(x,t,\nabla u) = \mu, \text{ where the right side is a general measure, } b \text{ is a strictly increasing } C^1\text{-function}, -\operatorname{div}(a(x,t,\nabla u)) \text{ is a Leray-Lions type operator with growth } |\nabla u|^{p-1} \text{ in } \nabla u \text{ and } H(x,t,\nabla u) \text{ is a non-linear lower order term which satisfy the growth condition with respect to } \nabla u.$ 

Key Words: Nonlinear parabolic equations, Existence, Renormalized solutions, Lower order term, Measure data.

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## 1. Introduction

Let  $\Omega$  be a bounded open subset of  $\mathbb{R}^N$ ,  $(N \ge 1)$ , T > 0 and let  $Q_T := \Omega \times (0,T)$ . We prove the existence of a renormalized solution for a class of nonlinear parabolic equations of the type:

$$\frac{\partial b(u)}{\partial t} - \operatorname{div}\left(a(x, t, \nabla u)\right) + H(x, t, \nabla u) = \mu \quad \text{in } Q_T, \tag{1.1}$$

$$b(u)(t=0) = b(u_0)$$
 in  $\Omega$ , (1.2)

$$u = 0 \quad \text{on } \partial\Omega \times (0, T). \tag{1.3}$$

In Problem (1.1)-(1.3) the framework is the following: the data  $\mu$  is a general measure, b is a strictly increasing  $C^1$ -function, the operator  $-\operatorname{div}(a(x,t,\nabla u))$  is a Leray–Lions operator which is coercive and grows like  $|\nabla u|^{p-1}$  with respect to  $\nabla u$  and  $H(x,t,\nabla u)$  is a nonlinear lower order. In the case where b(u) = u, H = 0, and

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the right hand side is a bounded measure, the existence of a distributional solution was proved in [3], but due the lack of regularity of solution, the distributional formulation is not strong enough to provide uniqueness (see [27] for a counter example in the elliptic case). To overcome this difficulty the notion of renormalized solutions firstly introduced by R.J. DiPerna and P.-L.Lions in [8] for the study of Boltzmann equation was adapted to parabolic equations and elliptic equations with  $L^1$  data. When  $\mu$  is measure data that does not charge the sets of zero p-capacity (the so called diffuse measure, see the definition in the section 2 below) a notion of renormalized solution for Problem (1.1)-(1.3) was introduced in [11] for b(u) = uand H = 0. Similar result was proved in [5] when H = 0. In [22] the existence of renormalized solution was proved in the case where b(u) = u and H = 0 and  $\mu$  is a general measure (see also [6]). In [10] a similar notion of entropy solution is also defined and proved to be equivalent to the renormalized one. In this paper we use a new definition of renormalized solution which is adopted in [25] and [24] for the study of parabolic p-Laplacian equations with general measure data.

The paper is organized as follows. In section 2 we give some preliminaries on the concept of p-capacity. Section 3 will be devoted to set our main assumptions and definition of renormalized solution and the statement of the existence result, while in Section 4 we give the proof of our main result.

#### 2. Preliminaries on parabolic capacity

We introduce the notion of p-capacity associated to our problem (for further details see [21], [11]). Let  $Q_T = \Omega \times (0,T)$  for any fixed T > 0 and  $1 , and let us recall that <math>V = W_0^{1,p}(\Omega) \cap L^2(\Omega)$ , endowed with its natural norm  $\|.\|_{W_0^{1,p}(\Omega)} + \|.\|_{L^2(\Omega)}$  and

$$W = \left\{ u \in L^{p}(0,T;V), u_{t} \in L^{p'}(0,T;V') \right\},\$$

endowed with its natural norm  $\|.\|_{L^p(0,T;V)} + \|.\|_{L^{p'}(0,T;V')}$ , remark that W is continuously embedded in  $C([0,T], L^2(\Omega))$ , and if  $1 , then <math>\mathcal{C}_c^{\infty}(Q_T)$  is dense in W. Let  $U \subseteq Q_T$  be an open set, we define the parabolic p-capacity of U as

$$cap_p(U) = \inf \left\{ \|u\|_W : u \in W, u \ge \chi_U \text{ a.e. in } Q_T \right\}$$

where as usual we set  $inf\{\emptyset\} = +\infty$ , then for any Borel set  $B \subseteq Q_T$  we define

$$cap_p(B) = \inf \left\{ cap_p(U) : U \text{ open set of } Q_T, B \subseteq U \right\}.$$

We will denote by  $\mathcal{M}(Q_T)$  the set of all Radon measures with bounded variation on  $Q_T$ , while, as we already mentioned,  $\mathcal{M}_0(Q_T)$  the set of all measures with bounded total variation over  $Q_T$  that do not charge the sets of zero p-capacity, that is if  $\mu \in \mathcal{M}_0(Q_T)$ , then  $\mu(E) = 0$ , for all  $E \subseteq Q_T$  such that  $cap_p(E) = 0$ . In [11] the authors proved the following decomposition theorem: **Theorem 2.1.** Let  $\mu$  be a bounded measure on  $Q_T$ . If  $\mu \in \mathcal{M}_0(Q_T)$  then there exists  $(f, g_1, g_2)$  such that  $f \in L^1(Q_T)$ ,  $g_1 \in L^{p'}(0, T; W^{-1, p'}(\Omega))$ ,  $g_2 \in L^p(0, T; V)$  and

$$\int_{Q_T} \phi \ d\mu = \int_{Q_T} f\phi \ dx \ dt + \int_0^T \langle g_1, \phi \rangle \ dt - \int_0^T \langle \phi_t, g_2 \rangle \ dt \ \phi \in \mathcal{C}^\infty_c(Q_T).$$

Such a triplet  $(f, g_1, g_2)$  will called decomposition of  $\mu$ .

**Definition 2.2.** A sequence of measures  $(\mu_n)$  in  $Q_T$  is equidiffuse if for every  $\varepsilon > 0$  there exists  $\eta > 0$  such that for every Borel set  $E \subseteq Q_T$ ,

$$cap_p(E) < \eta \Rightarrow |\mu_n|(E) < \varepsilon \ \forall n \ge 1.$$

Let  $\rho_n$  be a sequence of mollifiers on  $Q_T$ , the following result is proved in [25]

**Proposition 2.3.** If  $\mu \in \mathcal{M}_0(Q_T)$ , then the sequence  $\rho_n * \mu$  is equidiffuse.

If  $\mu \in \mathcal{M}(Q_T)$ , thanks to a well known decomposition result (see for instance [14]), we can split it into a sum (uniquely determined) of its absolutely continuous part  $\mu_d$  with respect to p-capacity and its singular part  $\mu_s$ , that is  $\mu_s$  is concentrated on a set E of zero p-capacity. Hence, if  $\mu \in \mathcal{M}(Q_T)$ , we have

$$\mu = \mu_d + \mu_s = \mu_d + \mu_s^+ - \mu_s^-$$

## 3. Assumptions and definition of a renormalized solution

Throughout the paper, we assume that the following assumptions hold true:  $\Omega$  is a bounded open set on  $\mathbb{R}^N$   $(N \ge 1)$ , T > 0 is given and we set  $Q_T = \Omega \times (0, T)$ 

$$b: \mathbb{R} \to \mathbb{R} \text{ and}$$
 (3.1)

is a strictly increasing  $C^1$ -function with b(0)=0, and there exists  $\gamma>0$  and  $\Lambda>0$  such that

$$\gamma \le b'(s) \le \Lambda, \ \forall s \in \mathbb{R}.$$
(3.2)

$$a: Q_T \times \mathbb{R}^N \to \mathbb{R}^N$$
 is a Carathéodory function (3.3)

$$a(x,t,\xi).\xi \ge \alpha |\xi|^p,\tag{3.4}$$

for almost every  $(x,t) \in Q_T$ , for every  $\xi \in \mathbb{R}^N$ , where  $\alpha > 0$  is a given real number.

$$|a(x,t,\xi)| \le \beta(L(x,t) + |\xi|^{p-1}), \tag{3.5}$$

for almost every  $(x,t) \in Q_T$ , for every  $\xi \in \mathbb{R}^N$ , where  $\beta > 0$  is a given real number, L is a non negative function in  $L^{p'}(Q_T)$ .

$$[a(x,t,\xi) - a(x,t,\xi')][\xi - \xi'] > 0.$$
(3.6)

Let  $H: Q_T \times \mathbb{R}^N \to \mathbb{R}$  be Carathéodory function such that for a.e.  $(x,t) \in Q_T$ and for every  $\xi \in \mathbb{R}^N$ , the growth condition

$$|H(x,t,\xi)| \le g(x,t)|\xi|^{\delta}, \qquad (3.7)$$

is satisfied, with  $\delta = \frac{p(N+1)-N}{N+2}$  and g belongs to  $L^{N+2,1}(Q_T)$ .

$$\mu \in \mathcal{M}(Q_T),\tag{3.8}$$

$$u_0$$
 is an element of  $L^1(\Omega)$ . (3.9)

We use in the present paper the two Lorentz spaces  $L^{q,1}(Q_T)$  and  $L^{q,\infty}(Q_T)$ , see for example ([18], [19]) for references about Lorentz spaces  $L^{q,s}$ . if  $f^*$  denotes the decreasing rearrangement of a measurable function f,

$$f^*(r) = \inf \left\{ s \ge 0 : meas\{(x,t) \in Q_T : |f(x,t)| > s\} < r \right\}, \ r \in [0, meas(Q_T)],$$

 $L^{q,1}(Q_T)$  is the space of Lebesgue measurable functions such that

$$\|f\|_{L^{q,1}(Q_T)}^q = \int_0^{meas(Q_T)} f^* r^{\frac{1}{q}} \frac{dr}{r} < \infty$$

while  $L^{q,\infty}(Q)$  is the space of Lebesgue measurable functions such that

$$||f||_{L^{q,\infty}(Q_T)} = \sup_{r>0} r[meas\{(x,t) \in Q_T : |f(x,t)| > r\}]^{\frac{1}{q}} < \infty.$$

If  $1 < q < \infty$  we have the generalized Hölder inequality

A

$$f \in L^{q,\infty}(Q_T), \forall g \in L^{q',1}(Q_T) \text{ such that } \frac{1}{q} + \frac{1}{q'} = 1,$$
$$\int_Q |fg| \le \|f\|_{L^{q,\infty}(Q_T)} \|g\|_{L^{q',1}(Q_T)}.$$

Now we give the definition of renormalized solution of Problem (1.1) - (1.3).

**Definition 3.1.** A measurable function u is a renormalized solution of Problem (1.1)-(1.3) if

$$T_k(u) \in L^p(0,T; W_0^{1,p}(\Omega)), \text{ for every } k > 0, \quad H(x,t,\nabla u) \in L^1(Q_T), \quad (3.11)$$

and if there exists a sequence of measures  $\Gamma^k \in \mathcal{M}(Q_T)$  such that:

1

$$\Gamma^{\kappa} \to \mu_s \text{ tightly as } k \to \infty,$$
 (3.12)

(3.10)

$$B_k(u)_t - \operatorname{div}\left(a(x,t,\nabla T_k(u))\right) + H(x,t,\nabla T_k(u)) = \mu_d + \Gamma^k \text{ in } \mathcal{D}'(Q_T),$$
  
where  $B_k(s) = \int_0^s T'_k(r)b'(r) \, dr, \ \forall s \in \mathbb{R}.$ 

**Remark 3.2.** Observe that (3.2) and (3.11) imply that each term in (3.13) is well defined and that (3.13) implies that  $B_k(u)_t - \operatorname{div}\left(a(x,t,\nabla T_k(u))\right) + H(x,t,\nabla T_k(u))$  is a bounded measure, then we have

$$B_k(u)_t - \operatorname{div}\left(a(x, t, \nabla T_k(u))\right) + H(x, t, \nabla T_k(u)) = \mu_d + \Gamma^k \text{ in } \mathcal{M}(Q_T).$$

A remark on the assumption (3.2) is also necessary. As one could check later, since the data is a measure  $\mu$ , we are forced to assume  $\gamma \leq b'(s) \leq \Lambda$ . We conjecture that this assumption is only technical and could be removed in order to deal with more general elliptic-parabolic problems (see for instance [1], [7]).

In order to prove the existence result we give the following Lemma

**Lemma 3.3.** Let u be a measurable function satisfying  $T_k(u) \in L^p(0, T; W_0^{1,p}(\Omega)) \cap L^{\infty}(0, T; L^2(\Omega))$  for every k > 0 such that:

$$\sup_{t\in(0,T)}\int_{\Omega}|T_k(u)|^2\,dx+\int_{Q_T}|\nabla T_k(u)|^p\,dxdt\leq Mk\ \, \forall k>0,$$

where M is a positive constant. Then

$$\begin{aligned} ||u|^{p-1}||_{L^{\frac{p(N+1)-N}{N(p-1)},\infty}(Q_T)} &\leq CM^{(\frac{p}{N}+1)\frac{N}{N+p'}}|Q_T|^{\frac{1}{p'}\frac{N}{N+p'}},\\ ||\nabla u|^{p-1}||_{L^{\frac{p(N+1)-N}{(N+1)(p-1)},\infty}(Q_T)} &\leq CM^{\frac{(N+2)(p-1)}{p(N+1)-N}}, \end{aligned}$$

where C is a constant which depends only on N and p.

*Proof.* See [13] and [12].

#### 4. Existence result

Let us introduce the following regularization of the data: for  $n \ge 1$  fixed

$$u_0^n \in C_c^{\infty}(Q_T)$$
, such that  $u_0^n \to u_0$  in  $L^1(\Omega)$ , (4.1)

$$\mu^n \in C^{\infty}(Q_T), \ \mu^n = \mu^n_d + \mu^n_s, \tag{4.2}$$

where  $\mu_d^n = \rho_n * \mu_d$  and  $\mu_s^n = \rho_n * \mu_s^+ - \rho_n * \mu_s^- = \lambda_+^n - \lambda_-^n$ . Moreover we have

$$\|\mu^n\|_{L^1(Q_T)} \le |\mu|_{\mathcal{M}(Q_T)},$$

and

 $\mu^n$  converges to  $\mu$  in the narrow topology of measures.

Let us now consider the following regularized problem

$$u^n \in L^p(0,T; W^{1,p}_0(\Omega)),$$
 (4.3)

$$\int_{0}^{T} \langle \frac{\partial v^{n}}{\partial t}, \varphi \rangle dt + \int_{Q_{T}} a(x, t, \nabla u^{n}) \nabla \varphi \, dx dt + \int_{Q_{T}} H(x, t, \nabla u^{n}) \varphi \, dx dt = \int_{Q_{T}} \mu^{n} \varphi \, dx dt$$
$$\forall \varphi \in L^{p}(0, T; W_{0}^{1, p}(\Omega)) \cap L^{\infty}(Q_{T}),$$
$$b(u^{n})(t=0) = b(u_{0}^{n}) \text{ in } \Omega, \qquad (4.4)$$

where  $v^n = b(u^n)$ .

As a consequence, proving existence of a weak solution  $u^n \in L^p(0,T; W_0^{1,p}(\Omega))$  of (4.3)-(4.4) is classical (see for instance [15]).

Now we give the following proposition which gives some compactness results.

**Proposition 4.1.** Let  $u^n$  and  $v^n$  be defined as before. Then

$$\||\nabla u^{n}|^{\delta}\|_{L^{\frac{N+2}{N+1},\infty}(Q_{T})} \le C,$$
(4.5)

$$\|u^n\|_{L^{\infty}(0,T;L^1(\Omega))} \le C, (4.6)$$

$$\int_{Q} |\nabla T_k(u^n)|^p \, dx dt \le Ck,\tag{4.7}$$

$$\int_{Q} |\nabla T_k(v^n)|^p \, dx dt \le Ck,\tag{4.8}$$

$$u^n \text{ is bounded in } L^q(0,T; W^{1,q}_0(\Omega)) \quad \forall \ 1 < q < p - \frac{N}{N+1},$$
 (4.9)

Moreover, there exists a measurable function u and v = b(u) such that  $T_k(u)$  and  $T_k(v)$  belong to  $L^p(0,T; W_0^{1,p}(\Omega))$ , and u belongs to  $L^{\infty}(0,T; L^1(\Omega))$ , up to a subsequence, for any k > 0 and for any  $1 < q < p - \frac{N}{N+1}$  we have

(4.10)

(4.13)

 $u^n \rightarrow u \text{ a.e. on } Q_T \text{ weakly in } L^q(0,T;W^{1,q}_0(\Omega)) \text{ and strongly in } L^1(Q_T),$ 

$$T_k(u^n) \rightharpoonup T_k(u)$$
 weakly in  $L^p(0,T; W_0^{1,p}(\Omega))$  and a.e. in  $Q_T$ , (4.11)

$$T_k(v^n) \rightharpoonup T_k(v)$$
 weakly in  $L^p(0,T; W_0^{1,p}(\Omega))$  and a.e. in  $Q_T$ . (4.12)

*Proof.* The proof of this Proposition relies on standard techniques for problems of type (4.3)-(4.4). Let k > 0, we take  $T_k(u^n)\chi_{(0,t)}$  as test function in (4.3) for every  $t \in (0,T)$  and we have

$$\int_{\Omega} \overline{B}_k(u^n)(t) \, dx + \int_{Q_t} a(x, t, \nabla u^n) \nabla T_k(u^n) \, dx dt$$
$$\leq \int_{Q_t} |H(x, t, \nabla u^n)| |T_k(u^n)| \, dx dt + \int_{Q_t} \mu^n T_k(u^n) \, dx dt + \int_{\Omega} \overline{B}_k(u^n_0) \, dx,$$

where  $\overline{B}_k(s) = \int_0^s T_k(r)b'(r) dr$ . Using (3.4) and (3.7) we obtain

$$\int_{\Omega} \overline{B}_k(u^n)(t) \, dx + \alpha \int_{Q_t} |\nabla T_k(u^n)|^p \, dx dt$$
$$\leq k \Big( \int_{Q_t} |g(x,t)| |\nabla u^n|^\delta \, dx dt + \|\mu^n\|_{L^1(Q_t)} + \|b(u_0^n)\|_{L^1(\Omega)} \Big)$$

if we take the supremum for  $t \in (0, t_1)$ , where  $t_1 \in (0, T)$  will be choosen later, by (3.2) we have

$$\frac{\gamma}{2} \sup_{t \in (0,t_1)} \int_{\Omega} |T_k(u^n)|^2 \, dx + \alpha \int_{Q_{t_1}} |\nabla T_k(u^n)|^p \, dx dt$$
$$\leq k \Big( \int_{Q_{t_1}} |g(x,t)| |\nabla u^n|^\delta \, dx dt + \|\mu^n\|_{L^1(Q_T)} + \|b(u_0^n)\|_{L^1(\Omega)} \Big),$$

and thanks to the generalized Hölder inequality we obtain

$$\frac{\gamma}{2} \sup_{t \in (0,t_1)} \int_{\Omega} |T_k(u^n)|^2 \, dx + \alpha \int_{Q_{t_1}} |\nabla T_k(u^n)|^p \, dx dt \tag{4.14}$$

$$\leq k \Big( \||\nabla u^n|^{\delta}\|_{L^{\frac{N+2}{N+1},\infty}(Q_{t_1})} \|g\|_{L^{N+2,1}(Q_{t_1})} + \|\mu^n\|_{L^1(Q_T)} + \|b(u_0^n)\|_{L^1(\Omega)} \Big) \leq Mk,$$

where  $M = \||\nabla u^n|^{\delta}\|_{L^{\frac{N+2}{N+1},\infty}(Q_{t_1})} \|g\|_{L^{N+2,1}(Q_{t_1})} + \|\mu^n\|_{L^1(Q_T)} + \|b(u_0^n)\|_{L^1(\Omega)},$  by Lemma 3.3 we obtain

$$\||\nabla u^{n}|^{\delta}\|_{L^{\frac{N+2}{N+1},\infty}(Q_{t_{1}})} = \||\nabla u^{n}|^{p-1}\|_{L^{\frac{p(N+1)-N}{(N+1)(p-1)},\infty}(Q_{t_{1}})}^{\frac{\delta}{p-1}}$$
(4.15)

$$\leq C(\||\nabla u^n|^{\delta}\|_{L^{\frac{N+2}{N+1},\infty}(Q_{t_1})} \|g\|_{L^{N+2,1}(Q_{t_1})} + \|\mu^n\|_{L^1(Q_T)} + \|b(u_0^n)\|_{L^1(\Omega)}).$$

If we choose  $t_1$  such that

$$1 - C \|g\|_{L^{N+2,1}(Q_{t_1})} > 0, (4.16)$$

holds, then we have

$$\||\nabla u^{n}|^{\delta}\|_{L^{\frac{N+2}{N+1},\infty}(Q_{t_{1}})} \le C,$$
(4.17)

which yields (4.5). Since  $\overline{B}_k(s) \ge \gamma \int_0^s T_1(r) dr \ge \gamma(|s|-1) \ \forall s \in \mathbb{R}$ , we obtain $\|u^n\|_{L^{\infty}(0,t_1;L^1(\Omega))} \le \frac{1}{\gamma}M + \ meas(\Omega).$  From (4.17) it follows that

$$\|u^n\|_{L^{\infty}(0,t_1;L^1(\Omega))} \le C. \tag{4.18}$$

Now we use the same technique as in ([23]). We consider a partition of the interval [0, T] into a finite number of intervals  $[0, t_1]$ ,  $[t_1, t_2], ..., [t_{n-1}, T]$  such that for each  $[t_{i-1}, t_i]$  the condition (4.16) holds.

In this way in each cylindre  $\Omega \times [t_{i-1}, t_i]$  we obtain a priori estimates of type (4.5) and (4.6). From (4.14) and (4.17) with T in place of  $t_1$  we obtain (4.7).

By using (4.6) and (4.7), and thanks to L. Boccardo and T. Gallouët (see [3]) we obtain (4.9). By (3.2), (4.9), and since  $\mu^n$  is bounded in  $L^1(Q_T)$ , one obtain that  $\frac{\partial v^n}{\partial t}$  is bounded in  $L^1(0,T;W^{-1,q'}(\Omega))$  for every  $q' < 1 + \frac{1}{(p-1)(N+1)}$ , using a standard compactness arguments (see [26]) yield (4.10), (4.11) and (4.12).

Let us introduce for  $k \geq 0$  fixed, the time regularization of the function  $T_k(v)$ . This kind of regularization has been first introduced by R. Landes. More recently, it has been exploited to solve a few nonlinear evolution problems with  $L^1$  or measure data. This specific time regularization of  $T_k(v)$  (for fixed  $k \geq 0$ ) is defined as follows. Let  $(v_0^{\nu})_{\nu}$  in  $L^{\infty}(\Omega) \cap W_0^{1,p}(\Omega)$  such that  $\|v_0^{\nu}\|_{L^{\infty}(\Omega)} \leq k$ , for all  $\nu > 0$ , and  $v_0^{\nu} \to T_k(b(u_0))$  a.e. in  $\Omega$  with  $\frac{1}{\nu} \|v_0^{\nu}\|_{L^p(\Omega)} \to 0$  as  $\nu \to +\infty$ . For fixed  $k \geq 0$  and  $\nu > 0$ , let us consider the unique solution  $T_k(v)_{\nu} \in L^{\infty}(Q_T) \cap V_0^{1,p}(v_0^{\nu})$ 

 $L^p(0,T,W_0^{1,\overline{p}}(\Omega))$  of the monotone problem:

$$\frac{\partial T_k(v)_{\nu}}{\partial t} + \nu (T_k(v)_{\nu} - T_k(v)) = 0 \text{ in } \mathcal{D}'(Q_T)$$
$$T_k(v)_{\nu}(t=0) = v_0^{\nu} \text{ in } \Omega.$$

The behavior of  $T_k(v)_{\nu}$  as  $\nu \to +\infty$  is investigated in [17] and we just recall here that:

$$T_k(v)_{\nu} \to T_k(v)$$
 strongly in  $L^p(0,T,W_0^{1,p}(\Omega))$  a.e. in  $Q_T$  as  $\nu \to +\infty$ 

with  $||T_k(v)_{\nu}||_{L^{\infty}(\Omega)} \leq k$  for any  $\nu > 0$ , and  $\frac{\partial T_k(v)_{\nu}}{\partial t} \in L^p(0, T, W_0^{1, p}(\Omega))$ . We will denote  $\omega(n, \nu, k, \varepsilon)$  any quantity that vanishes as the parameters go to their limit point with in the same order in which they appear, that is, for example

$$\overline{\lim_{\varepsilon \to 0}} \ \overline{\lim_{k \to \infty}} \ \overline{\lim_{\nu \to \infty}} \ \overline{\lim_{n \to \infty}} \ |\omega(n, \nu, k, \varepsilon)| = 0.$$

We give the following result which has been proved in [2].

**Lemma 4.2.** Let  $v^n$  be a sequence in  $L^p(0,T; W_0^{1,p}(\Omega)) \cap C^0([0,T]; L^2(\Omega))$ , and  $(v^n)_t \in L^{p'}(0,T; W^{-1,p'}(\Omega))$ , suppose that  $v^n$  converges almost everywhere in  $Q_T$  to a function v such that  $T_k(v) \in L^p(0,T; W_0^{1,p}(\Omega))$  for every k > 0. then we have

$$\int_0^T \langle \frac{\partial v^n}{\partial t}, T_{\varepsilon}(v^n - T_k(v)_{\nu})) \rangle \, dt \ge \omega(n, \nu, k, \varepsilon).$$

# **Proposition 4.3.** The sequence $(\nabla u^n)$ converges to $\nabla u$ a.e. in $Q_T$ .

*Proof.* Adopting the method used in [2], we prove that for some  $\theta > 0$ , one has up to subsequences still denoted by  $u^n$  (for simplicity of notation, we will omit the dependence of a on x and t),

$$\left[ (a(\nabla u^n) - a(\nabla u)).(\nabla u^n - \nabla u) \right]^{\theta} \to 0 \text{ a.e. in } Q.$$
(4.19)

Note that (4.19) will be true if we show that

$$\int_{Q_T} \left[ (a(\nabla u^n) - a(\nabla u)) . (\nabla u^n - \nabla u) \right]^{\theta} dx dt = \omega(n)$$
(4.20)

The same argument in [16] and under asymptions on  $a(x, t, \xi)$  implies that  $\nabla u^n$  converges to  $\nabla u$  a.e. in  $Q_T$ .

Thanks to Proposition 4.1, the following estimate holds

$$meas(\{|v| \ge k\}) = \omega(k),$$

We can write

$$\int_{Q_T} \left[ (a(\nabla u^n) - a(\nabla u)).(\nabla u^n - \nabla u) \right]^{\theta} dx dt$$
  
= 
$$\int_{\{|v| \ge k\}} \left[ (a(\nabla u^n) - a(\nabla u)).(\nabla u^n - \nabla u) \right]^{\theta} dx dt$$
  
+ 
$$\int_{\{|v| < k\}} \left[ (a(\nabla u^n) - a(\nabla u)).(\nabla u^n - \nabla u) \right]^{\theta} dx dt$$
  
= 
$$I_{n,k} + J_{n,k}.$$

Since  $u^n$  is bounded in  $L^q(0,T;W^{1,q}_0(\Omega))$  for  $q , we can choose <math>\theta < \frac{q}{p} < 1$ , so that using Hölder inequality, we obtain

$$|I_{n,k}| \le c \ meas(\{|v| \ge k\})^{1-\theta p/q},$$

and then  $I_{n,k} = \omega(k)$ . Now we set

$$\Psi_{n,k} = \left(a(\nabla u^n) - a(\nabla u\chi_{\{|v| < k\}})\right) \left(\nabla u^n - \nabla u\chi_{\{|v| < k\}}\right),$$

and we have

$$\int_{Q_T} \left[ (a(\nabla u^n) - a(\nabla u)) . (\nabla u^n - \nabla u) \right]^{\theta} dx dt$$

$$\leq \int_{Q_T} \Psi^{\theta}_{n,k} \chi_{\{|v^n - T_k(v)_{\nu}| \le \varepsilon\}} + \int_{Q_T} \Psi^{\theta}_{n,k} \chi_{\{|v^n - T_k(v)_{\nu}| > \varepsilon\}} + \omega(k),$$
(4.21)

since  $\Psi_{n,k}^{\theta}$  is bounded in  $L^{q/\theta p}(Q_T)$  independently of n and k,  $\chi_{\{|v^n - T_k(v)_\nu| > \varepsilon\}}$  converges to  $\chi_{\{|v - T_k(v)| > \varepsilon\}}$  almost everywhere in  $Q_T$  as n tends to  $+\infty$  (see [2],

Lemma 3.2 ) and  $\chi_{\{|v-T_k(v)_\nu|>\varepsilon\}}$  converges to zero almost everywhere in  $Q_T$  as  $\nu$  and k tends to  $+\infty$  we obtain

$$\int_{Q_T} \Psi^{\theta}_{n,k} \chi_{\{|v^n - T_k(v)_\nu| > \varepsilon\}} = \omega(n,\nu,k),$$

using Hölder inequality, (4.21) becomes

$$\int_{Q_T} \left[ (a(\nabla u^n) - a(\nabla u)) . (\nabla u^n - \nabla u) \right]^{\theta} dx dt$$
  
$$\leq meas(Q_T)^{1-\theta} \left( \int_{Q_T} \Psi_{n,k} \chi_{\{|v^n - T_k(v)_\nu| \leq \varepsilon\}} \right)^{\theta} + \omega(n,\nu,k).$$

Then it remains to prove that

$$\int_{Q_T} \Psi_{n,k} \chi_{\{|v^n - T_k(v)_\nu| \le \varepsilon\}} = \omega(n,\nu,k,\varepsilon).$$
(4.22)

By assumption (3.2) we can write

$$\int_{Q_T} \Psi_{n,k} \chi_{\{|v^n - T_k(v)_\nu| \le \varepsilon\}}$$

$$\leq \frac{1}{\gamma} \Big( \int_{Q_T} b'(u^n) a(\nabla u^n) \Big( \nabla u^n - \nabla u \chi_{\{|v| \le k\}} \Big) \chi_{\{|v^n - T_k(v)_\nu| \le \varepsilon\}} \Big)$$

$$- \frac{1}{\gamma} \Big( \int_{Q_T} b'(u^n) a(\nabla u \chi_{\{|v| \le k\}}) \Big( \nabla u^n - \nabla u \chi_{\{|v| \le k\}} \Big) \chi_{\{|v^n - T_k(v)_\nu| \le \varepsilon\}} \Big)$$

$$= \frac{1}{\gamma} \Big( \int_{Q_T} b'(u^n) a(\nabla u \chi_{\{|v| \le k\}}) \Big( \nabla u^n - \nabla u \chi_{\{|v| \le k\}} \Big) \chi_{\{|v^n - T_k(v)_\nu| \le \varepsilon\}} \Big)$$

By Proposition 4.1 and since  $|T_k(v)_{\nu}| \leq k$  we obtain

$$\int_{Q_T} b'(u^n) a(\nabla u \chi_{\{|v| \le k\}}) \Big( \nabla u^n - \nabla u \chi_{\{|v| \le k\}} \Big) \chi_{\{|v^n - T_k(v)_\nu| \le \varepsilon\}}$$
(4.24)  

$$= \int_{Q_T} a(\nabla u \chi_{\{|v| \le k\}}) \Big( \nabla v^n - b'(u^n) b'(u)^{-1} \nabla v \chi_{\{|v| \le k\}} \Big) \chi_{\{|v^n - T_k(v)_\nu| \le \varepsilon\}}$$
  

$$= \int_{Q_T} a(\nabla u \chi_{\{|v| \le k\}}) \Big( \nabla v^n - \nabla T_k(v)_\nu \Big) \chi_{\{|v^n - T_k(v)_\nu| \le \varepsilon\}}$$
  

$$+ \int_{Q_T} a(\nabla u \chi_{\{|v| \le k\}}) \Big( \nabla T_k(v)_\nu - b'(u^n) b'(u)^{-1} \nabla T_k(v) \Big) \chi_{\{|v^n - T_k(v)_\nu| \le \varepsilon\}}$$
  

$$= A_1 + A_2.$$

For  $\varepsilon < 1$  and thanks to Propsition 4.1 we obtain

$$A_1 = \int_{Q_T} a(\nabla u \chi_{\{|v| \le k\}}) \nabla T_{\varepsilon}(v^n - T_k(v)_{\nu})$$

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$$= \int_{Q_T} a(\nabla u \chi_{\{|v| \le k\}}) \nabla T_{\varepsilon}(T_{k+1}(v^n) - T_k(v)_{\nu})$$
$$= \int_{Q_T} a(\nabla u \chi_{\{|v| \le k\}}) \nabla T_{\varepsilon}(T_{k+1}(v) - T_k(v)_{\nu}) + \omega(n)$$

and the strong convergence of  $\nabla T_k(v)_{\nu}$  to  $\nabla T_k(v)$  in  $(L^p(Q_T))^N$  leads to

$$A_1 = \int_{Q_T} a(\nabla u \chi_{\{|v| \le k\}}) \nabla T_{\varepsilon}(T_{k+1}(v) - T_k(v)) + \omega(n, \nu)$$
$$= \omega(n, \nu).$$

By Proposition 4.1 we have  $b'(u^n)$  converges to b'(u) almost everywhere in  $Q_T$ , since  $a(\nabla u\chi_{\{|v|\leq k\}})$  belongs to  $(L^{p'}(Q_T))^N$ ,  $\nabla T_k(v)_{\nu}$  and  $\nabla T_k(v)$  belong to  $(L^p(Q_T))^N$ , the Lebesgue's convergence theorem leads to

$$|A_{2}| \leq \int_{Q_{T}} |a(\nabla u\chi_{\{|v| \leq k\}})| |\nabla T_{k}(v)_{\nu} - b'(u^{n})b'(u)^{-1}\nabla T_{k}(v)|,$$
  
$$\leq \int_{Q_{T}} |a(\nabla u\chi_{\{|v| \leq k\}})| |\nabla T_{k}(v)_{\nu} - \nabla T_{k}(v)| + \omega(n),$$

and by the strong convergence of  $\nabla T_k(v)_{\nu}$  to  $\nabla T_k(v)$  in  $(L^p(Q_T))^N$  we obtain

$$A_2 = \omega(n, \nu).$$

On the other hand we have

$$\int_{Q_T} b'(u^n) a(\nabla u^n) \Big( \nabla u^n - \nabla u \chi_{\{|v| \le k\}} \Big) \chi_{\{|v^n - T_k(v)_\nu| \le \varepsilon\}}$$

$$= \int_{Q_T} a(\nabla u^n) \nabla (v^n - T_k(v)_\nu) \chi_{\{|v^n - T_k(v)_\nu| \le \varepsilon\}}$$

$$+ \int_{Q_T} a(\nabla u^n) \Big( \nabla T_k(v)_\nu - b'(u^n) b'(u)^{-1} \nabla T_k(v) \Big) \chi_{\{|v^n - T_k(v)_\nu| \le \varepsilon\}}.$$
(4.25)

We deal with the second term on the right side of (4.25), by assumption (3.1) it is clear that  $\{|v^n| \leq k + \varepsilon\} \subset \{|u^n| \leq k_{\varepsilon} = \max\{b^{-1}(k + \varepsilon), |b^{-1}(-k - \varepsilon)|\}\}$  and by Hölder inequality we have

$$\left| \int_{Q_T} a(\nabla u^n) \Big( \nabla T_k(v)_{\nu} - b'(u^n) b'(u)^{-1} \nabla T_k(v) \Big) \chi_{\{|v^n - T_k(v)_{\nu}| \le \varepsilon\}} \right|$$
  
$$\leq \|a(\nabla T_{k_{\varepsilon}}(u^n)\|_{L^{p'}(Q_T)} \|\nabla T_k(v)_{\nu} - b'(u^n) b'(u)^{-1} \nabla T_k(v)\|_{L^p(Q_T)},$$

the almost everywhere convergence of  $b'(u^n)$  to b'(u) and Lebesgue's convergence theorem imply that  $b'(u^n)b'(u)^{-1}\nabla T_k(v)$  converges to  $\nabla T_k(v)$  strongly in  $(L^p(Q_T))^N$ , since  $|a(\nabla T_{k_{\varepsilon}}(u^n)|$  is bounded in  $L^{p'}(Q_T)$  we obtain

$$\left|\int_{Q_T} a(\nabla u^n) \Big(\nabla T_k(v)_{\nu} - b'(u^n)b'(u)^{-1} \nabla T_k(v)\Big) \chi_{\{|v^n - T_k(v)_{\nu}| \le \varepsilon\}}\right|$$

$$\leq C \|\nabla T_k(v)_{\nu} - \nabla T_k(v)\|_{L^p(Q_T)} + \omega(n),$$

and the strong convergence of  $\nabla T_k(v)_{\nu}$  to  $\nabla T_k(v)$  in  $(L^p(Q_T))^N$  leads to

$$\int_{Q_T} b'(u^n) a(\nabla u^n) \Big( b'(u^n)^{-1} \nabla T_k(v)_{\nu} - b'(u)^{-1} \nabla T_k(v) \Big) \chi_{\{|v^n - T_k(v)_{\nu}| \le \varepsilon\}} = \omega(n,\nu).$$

Hence (4.23), (4.24) and (4.25) imply that

$$\int_{Q_T} \Psi_{n,k} \chi_{\{|v^n - T_k(v)_\nu| \le \varepsilon\}} \le \int_{Q_T} a(\nabla u^n) \nabla (v^n - T_k(v)_\nu) \chi_{\{|v^n - T_k(v)_\nu| \le \varepsilon\}} + \omega(n,\nu).$$

Now we use the equation solved by  $u^n$ . Taking  $T_{\varepsilon}(v^n - T_k(v)_{\nu})$  in (4.3) we obtain

$$\int_0^T \langle \frac{\partial v^n}{\partial t}, T_{\varepsilon}(v^n - T_k(v)_{\nu}) \rangle dt + \int_{Q_T} a(\nabla u^n) \nabla T_{\varepsilon}(v^n - T_k(v)_{\nu}) dx dt + \int_{Q_T} H(x, t, \nabla u^n) T_{\varepsilon}(v^n - T_k(v)_{\nu}) dx dt = \int_{Q_T} \mu^n T_{\varepsilon}(v^n - T_k(v)_{\nu}) dx dt.$$

By property of  $\mu^n$  we have

$$\left|\int_{Q_T} \mu^n T_{\varepsilon}(v^n - T_k(v)_{\nu}) \, dx dt\right| \le \varepsilon \|\mu^n\|_{L^1(Q_T)} \le \varepsilon |\mu|_{\mathcal{M}(Q_T)}$$

By generalized Hölder inequality we have

$$\left| \int_{Q_T} H(x, t, \nabla u^n) T_{\varepsilon}(v^n - T_k(v)_{\nu}) \, dx dt \right| \le \varepsilon \|g\|_{L^{N+2,1}(Q_T)} \||\nabla u^n|^{\delta}\|_{L^{\frac{N+2}{N+1},\infty}(Q_T)}.$$

By Lemma 4.2 we obtain

$$\int_{Q_T} a(\nabla u^n) \nabla T_{\varepsilon}(v^n - T_k(v)_{\nu}) \, dx dt \le \varepsilon \Big( C \|g\|_{L^{N+2,1}(Q_T)} + |\mu|_{\mathcal{M}(Q_T)} \Big).$$

Hence

$$\int_{Q_T} a(\nabla u^n) \nabla T_{\varepsilon}(v^n - T_k(v)_{\nu}) \, dx dt \le \omega(n, \nu, \varepsilon).$$
(4.26)

Then by (4.26) we obtain (4.22) and therefore (4.20) and (4.19).

**Remark 4.4.** Let us observe that from Proposition 4.3 we have  $H(x, t, \nabla u^n)$  converges to  $H(x, t, \nabla u)$  a.e. in  $Q_T$  and by Proposition 4.1  $H(x, t, \nabla u^n)$  is equiintegrable in  $L^1(Q_T)$ . Indeed if E is a measurable set of  $Q_T$ , due the growth assumption (3.7) on H, estimate (4.5) yields that

$$\int_{E} |H(x,t,\nabla u^{n})| \, dxdt \leq \int_{E} g(x,t) |\nabla u^{n}|^{\delta} \, dxdt$$

$$\leq C \|g\|_{L^{N+2,1}(E)}.$$

We conclude that  $H(x,t,\nabla u^n)$  is equi-integrable in  $L^1(Q_T)$ . Then by Vitali's theorem we deduce that  $H(x,t,\nabla u^n)$  converges to  $H(x,t,\nabla u)$  strongly in  $L^1(Q_T)$ . Let also remark that from Proposition 4.3, assumption (3.5) on a and Vitali's theorem, we deduce that  $a(x,t,\nabla u^n)$  is strongly compact in  $L^1(Q_T)$ .

Now we define the space S by

$$\mathcal{S} = \Big\{ z \in L^p(0,T; W_0^{1,p}(\Omega)), z_t \in L^{p'}(0,T; W^{-1,p'}(\Omega)) + L^1(Q_T) \Big\},\$$

endowed with its natural norm  $\|.\|_{L^{p}(0,T;W_{0}^{1,p}(\Omega))} + \|.\|_{L^{p'}(0,T;W^{-1,p'}(\Omega))+L^{1}(Q_{T})}$ , and its sub-space  $\mathcal{W}_{1}$  as

$$\mathcal{W}_1 = \Big\{ z \in L^p(0,T; W_0^{1,p}(\Omega)) \cap L^\infty(Q_T), z_t \in L^{p'}(0,T; W^{-1,p'}(\Omega)) + L^1(Q_T) \Big\},\$$

endowed with its natural norm

$$\|.\|_{L^{p}(0,T;W_{0}^{1,p}(\Omega))} + \|.\|_{L^{\infty}(Q_{T})} + \|.\|_{L^{p'}(0,T;W^{-1,p'}(\Omega)) + L^{1}(Q_{T})},$$

for any p > 1.

Let us recall that a function z is called  $cap_p$ -quasi continuous if for every  $\varepsilon > 0$ there exists an open set  $F_{\varepsilon}$  with  $cap(F_{\varepsilon}) \leq \varepsilon$  such that the restriction of z to  $Q_T \setminus F_{\varepsilon}$  is continuous. The following result shows that every functions in  $\mathcal{W}_1$ satisfy a capacitary estimate for the parabolic capacity.

**Theorem 4.5.** Let  $z \in W_1$ , then z admits a unique  $cap_p$ -quasi continuous representative. Moreover, we have

$$cap_p(\{|z| > k\}) \le \frac{C}{k}max\{[z]^{\frac{1}{p}}_{*}, [z]^{\frac{1}{p'}}_{*}\},\$$

where

$$[z]_* = \|z\|_{L^p(0,T;W_0^{1,p}(\Omega))}^p + \|z_t^1\|_{L^{p'}(0,T;W^{-1,p'}(\Omega))}^{p'} + \|z\|_{L^{\infty}(Q_T)}\|z_t^2\|_{L^1(Q_T)} + \|z\|_{L^{\infty}(0,T;L^2(\Omega))}^2,$$

such that  $z_t^1 \in L^{p'}(0,T; W^{-1,p'}(\Omega)), z_t^2 \in L^1(Q_T)$  is any decomposition of  $z_t$ , that is  $z_t = z_t^1 + z_t^2$ .

*Proof.* See [22], Theorem 3 and Lemma 2.

Now we prove the following theorem

**Theorem 4.6.** Let  $u^n \in L^p(0,T; W^{1,p}_0(\Omega))$  be a solution of Problem (4.3)-(4.4) then

$$cap_p(\{|v^n| > k\}) \le \frac{C}{k}max\{k^{\frac{1}{p}}, k^{\frac{1}{p'}}\}, \ \forall k \ge 1$$

*Proof.* Due to the presence of the lower order term H, the approach used in [25] in the proof of Theorem 1.2 does not apply here, to overcome this difficulty we are going to exploit the method used in [22] Theorem 4. Let us first introduce the following function

$$G_k(s) = \begin{cases} 1 & \text{if } |s| \le k, \\ k+1-|s| & \text{if } k < |s| \le k+1, \\ 0 & \text{if } |s| > k+1. \end{cases}$$

let us denote by  $\overline{G}_k(s)$  the primitive function of  $G_k(s)$ . Since we have

$$\int_{Q_T} |\nabla T_k(v^n)|^p \, dx dt \le Ck,$$

we obtain

$$\int_{Q_T} |\nabla \overline{G}_k(v^n)|^p \, dx dt \le Ck. \tag{4.27}$$

Given  $\varphi \in C_c^{\infty}(Q_T)$  and taking  $G_k(v^n)\varphi$  as test function in (4.3) we have in the sense of distribution

$$\overline{G}_k(v^n)_t = \operatorname{div}\Big(G_k(v^n)a(x,t,\nabla u^n)\Big)$$
(4.28)

$$\begin{aligned} -b'(u^{n})a(x,t,\nabla u^{n}).\nabla u^{n}\chi_{\{k\leq v^{n}< k+1\}} + b'(u^{n})a(x,t,\nabla u^{n}).\nabla u^{n}\chi_{\{-k-1< v^{n}\leq -k\}} \\ -H(x,t,\nabla u^{n})G_{k}(v^{n}) + G_{k}(v^{n})\mu_{n}, \end{aligned}$$

therefore by assumption (3.2) and Proposition 4.1, we have

$$\overline{G}_k(v^n)_t \in L^{p'}(0,T;W^{-1,p'}(\Omega)) + L^1(Q),$$

and

$$\overline{G}_k(v^n) \in L^p(0,T; W^{1,p}_0(\Omega)) \cap L^\infty(Q_T),$$

thanks to Theorem 4.5,  $\overline{G}_k(v^n)$  has a  $cap_p$ -quasicontinuous representative. To conclude the proof is enough to prove the capacitary estimate of  $v_n$ . Since  $\{|v^n| > k\} = \{\overline{G}_k(v^n) > k\}$ , by Theorem 4.5 we obtain

$$cap_p(\{|v^n| > k\}) \le \frac{C}{k}max\{[\overline{G}_k(v^n)]^{\frac{1}{p}}_*, [\overline{G}_k(v^n)]^{\frac{1}{p'}}_*\}.$$

Taking  $\theta_k(v^n) = T_{k+1}(v^n) - T_k(v^n)$  as test function in (4.3) leads to

$$\int_{\Omega} \Theta_k(v^n)(T) \, dx + \int_{\{k < |v^n| \le k+1\}} b'(u^n) a(x, t, \nabla u^n) \cdot \nabla u^n \, dx dt$$
$$+ \int_{Q_T} H(x, t, \nabla u^n) \theta_k(v^n) \, dx dt = \int_{Q_T} \theta_k(v^n) \mu^n \, dx dt + \int_{\Omega} \Theta_k(b(u_0^n)) \, dx,$$

where  $\Theta_k(s) = \int_0^s \theta_k(r) dr \ \forall s \in \mathbb{R}$ . Since  $\|\theta_k(v^n)\|_{L^{\infty}(Q_T)} \leq 1$  and  $H(x, t, \nabla u^n)$  is strongly compact in  $L^1(Q)$  one obtains

$$\int_{Q_T} b'(u^n) a(x,t,\nabla u^n) \cdot \nabla u^n \chi_{\{k \le v^n < k+1\}} \, dxdt \le C,$$
$$\int_{Q_T} b'(u^n) a(x,t,\nabla u^n) \cdot \nabla u^n \chi_{\{-k-1 < v^n \le -k\}} \, dxdt \le C,$$
$$\int_{Q_T} |H(x,t,\nabla u^n) G_k(v^n)| \, dxdt \le C.$$

Then, from (4.28) it follows that

$$\|\overline{G}_{k}(v^{n})_{t}^{1}\|_{L^{p'}(0,T;W^{-1,p'}(\Omega))}^{p'} \leq Ck,$$
$$\|\overline{G}_{k}(v^{n})_{t}^{2}\|_{L^{1}(Q)} \leq C,$$

using the following estimate

$$\|\overline{G}_{k}(v^{n})\|_{L^{\infty}(0,T;L^{2}(\Omega))}^{2} \leq \|\overline{G}_{k}(v^{n})\|_{L^{\infty}(Q_{T})}\|\overline{G}_{k}(v^{n})\|_{L^{\infty}(0,T;L^{1}(\Omega))},$$

since  $v^n$  is bounded in  $L^{\infty}(0,T;L^1(\Omega))$ , we conclude that

$$cap_p(\{|v^n| > k\}) \le \frac{C}{k}max\{k^{\frac{1}{p}}, k^{\frac{1}{p'}}\}.$$

We have the following technical result

**Lemma 4.7.** Let  $\mu_s = \mu_s^+ - \mu_s^- \in \mathcal{M}(Q_T)$  where  $\mu_s^+$  and  $\mu_s^-$  are concentrated respectively, on two disjoint  $E^+$  and  $E^-$  of zero p-capacity. Then, for every  $\delta > 0$ , there exist two compact sets  $K_{\delta}^+ \subseteq E^+$  and  $K_{\delta}^- \subseteq E^-$  such that

$$\mu_s^+(E^+ \backslash K_\delta^+) \le \delta, \quad \mu_s^-(E^+ \backslash K_\delta^-) \le \delta, \tag{4.29}$$

and there exist  $\psi_{\delta}^+$ ,  $\psi_{\delta}^- \in C_c^{\infty}(Q_T)$ , such that

$$\psi_{\delta}^+ \equiv 1 \text{ and } \psi_{\delta}^- \equiv 1 \text{ respectively on } K_{\delta}^+ \text{ and } K_{\delta}^-,$$
 (4.30)

$$0 \le \psi_{\delta}^+, \quad \psi_{\delta}^- \le 1, \tag{4.31}$$

$$supp(\psi_{\delta}^{+}) \cap supp(\psi_{\delta}^{-}) \equiv \emptyset.$$

$$(4.32)$$

Moreover

$$\|\psi_{\delta}^{+}\|_{\mathcal{S}} \le \delta, \quad \|\psi_{\delta}^{-}\|_{\mathcal{S}} \le \delta, \tag{4.33}$$

and in particular, there exists a decomposition of  $(\psi_{\delta}^+)_t$  and a decomposition of  $(\psi_{\delta}^-)_t$  such that

$$\|(\psi_{\delta}^{+})_{t}^{1}\|_{L^{p'}(0,T;W^{-1,p'}(\Omega))} \leq \frac{\delta}{3}, \quad \|(\psi_{\delta}^{+})_{t}^{2}\|_{L^{1}(Q_{T})} \leq \frac{\delta}{3}, \tag{4.34}$$

$$\|(\psi_{\delta}^{-})_{t}^{1}\|_{L^{p'}(0,T;W^{-1,p'}(\Omega))} \leq \frac{\delta}{3}, \quad \|(\psi_{\delta}^{-})_{t}^{2}\|_{L^{1}(Q_{T})} \leq \frac{\delta}{3}.$$
(4.35)

Both  $\psi_{\delta}^+$  and  $\psi_{\delta}^-$  converges to zero \*- weakly in  $L^{\infty}(Q_T)$ , in  $L^1(Q_T)$ , and up to subsequences, almost everywhere as  $\delta$  vanishes. Moreover, if  $\lambda_+^n$  and  $\lambda_-^n$  are as in (4.2) we have

$$\int_{Q_T} \psi_{\delta}^- d\lambda_+^n = \omega(n,\delta), \quad \int_{Q_T} \psi_{\delta}^- d\mu_s^+ \le \delta, \tag{4.36}$$

$$\int_{Q_T} \psi_{\delta}^+ d\lambda_-^n = \omega(n,\delta), \quad \int_{Q_T} \psi_{\delta}^+ d\mu_s^- \le \delta, \tag{4.37}$$

$$\int_{Q_T} (1 - \psi_{\delta}^+) \, d\lambda_+^n = \omega(n, \delta), \quad \int_{Q_T} (1 - \psi_{\delta}^+) \, d\mu_s^+ \le \delta, \tag{4.38}$$

$$\int_{Q_T} (1 - \psi_{\delta}^-) d\lambda_-^n = \omega(n, \delta), \quad \int_{Q_T} (1 - \psi_{\delta}^-) d\mu_s^- \le \delta.$$
Lemma 5.

Proof. See [22], Lemma 5.

Now we prove the following theorem

**Theorem 4.8.** Under assumptions (3.1)-(3.8), there exists at least a renormalized solution u of Problem (1.1)-(1.3).

Let us fix  $\sigma > 0$  and define

$$S_{k,\sigma}(s) = \begin{cases} 1 & \text{if } |s| \le k, \\ 0 & \text{if } |s| > k + \sigma, \\ \text{affine if otherwise.} \end{cases}$$

*Proof.* Step 1. Estimates in  $L^1(Q_T)$  on the energy term. Using  $h_{k,\sigma}(u^n) = \frac{1}{\sigma}(T_{k+\sigma}(u^n) - T_k(u^n))$  as test function in (4.3) we obtain

$$\int_{\Omega} B_{h_{k,\sigma}}^{\star}(u^{n})(T) dx + \frac{1}{\sigma} \int_{\{k < |u^{n}| \le k + \sigma\}} a(x,t,\nabla u^{n}) \cdot \nabla u^{n} dx dt$$
$$+ \int_{Q_{T}} H(x,t,\nabla u^{n}) h_{k,\sigma}(u^{n}) dx dt = \int_{Q_{T}} h_{k,\sigma}(u^{n}) \mu^{n} dx dt + \int_{\Omega} B_{h_{k,\sigma}}^{\star}(u_{0}^{n}) dx,$$

where  $B_{h_{k,\sigma}}^{\star}(s) = \int_{0}^{s} b'(r)h_{k,\sigma}(r) dr \ \forall s \in \mathbb{R}.$ So that dropping positive terms

$$\frac{1}{\sigma} \int_{\{k < |u^n| \le k + \sigma\}} a(x, t, \nabla u^n) \cdot \nabla u^n \, dx dt \tag{4.40}$$

$$\leq \int_{\{|u^n|>k\}} |\mu^n| \, dx dt + \int_{\{|u^n|>k\}} |H(x,t,\nabla u^n)| \, dx dt + \int_{\{|u^n_0|>k\}} b(u^n_0) \, dx.$$

which implies, in particular,

$$\frac{1}{\sigma} \int_{\{k < |u^n| \le k + \sigma\}} a(x, t, \nabla u^n) \cdot \nabla u^n \, dx dt \le C.$$
(4.41)

**Step 2.** Equation for the truncations. Given  $\varphi \in C_c^{\infty}(Q_T)$ , taking  $S_{k,\sigma}(u^n)\varphi$  as test function in (4.3), we obtain

$$B_{S_{k,\sigma}}^{\star}(u^{n})_{t} - \operatorname{div}\left(S_{k,\sigma}(u^{n})a(x,t,\nabla u^{n})\right) + H(x,t,\nabla u^{n})S_{k,\sigma}(u^{n})$$

$$= \mu_{d}^{n} + \mu_{s}^{n}S_{k,\sigma}(u^{n}) + \frac{1}{\sigma}sign(u^{n})a(x,t,\nabla u^{n}).\nabla u^{n}\chi_{\{k<|u^{n}|\leq k+\sigma\}}$$

$$+ \mu_{d}^{n}(S_{k,\sigma}(u^{n})-1) \text{ in } \mathcal{D}'(Q_{T}), \qquad (4.42)$$

where  $B^*_{S_{k,\sigma}}(s) = \int_0^s b'(r) S_{k,\sigma}(r) dr$ . From (4.41), there exists a bounded Radon measure  $\varsigma^n_k$  such that, as  $\sigma$  goes to zero

$$\frac{1}{\sigma}sign(u^n)a(x,t,\nabla u^n).\nabla u^n\chi_{\{k<|u^n|\leq k+\sigma\}}\rightharpoonup \varsigma_k^n \text{ * weakly in }\mathcal{M}(Q_T).$$

Taking the limit as  $\sigma$  vanishes in (4.42) it follows that

$$B_{k}(u^{n})_{t} - \operatorname{div}\left(a(x,t,\nabla T_{k}(u^{n}))\right) + H(x,t,\nabla T_{k}(u^{n})) = \mu_{d}^{n} + \mu_{s}^{n}\chi_{\{|u^{n}| \le k\}}$$
$$+\varsigma_{k}^{n} - \mu_{d}^{n}\chi_{\{|u^{n}| > k\}} \text{ in } \mathcal{D}'(Q_{T}),$$

where  $B_k(s) = \int_0^s T'_k(r)b'(r) dr$ ,  $\forall s \in \mathbb{R}$ . We define the measure  $\Gamma_n^k$  as

$$\Gamma_n^k = \mu_s^n \chi_{\{|u^n| \le k\}} + \varsigma_k^n - \mu_d^n \chi_{\{|u^n| \ge k\}}.$$

Notice that

$$\|\Gamma_n^k\|_{L^1(Q_T)} \le C,$$

so that there exist  $\Gamma^k \in \mathcal{M}(Q)$  such that

$$\Gamma_n^k \rightharpoonup \Gamma^k \star \text{weakly in } \mathcal{M}(Q_T).$$

Therefore, using Proposition 4.1 and Proposition 4.3, in the sense of distribution we have

(4.43)  
$$B_k(u)_t - \operatorname{div}\left(a(x, t, \nabla T_k(u))\right) + H(x, t, \nabla T_k(u^n)) = \mu_d + \Gamma^k \text{ in } \mathcal{D}'(Q_T).$$

Step 3. The limit of  $\Gamma^k$ . By subtracting (4.43) from the distributional formulation of (4.3) we obtain for any  $\varphi \in C_c^{\infty}(Q_T)$ 

$$\int_{Q_T} (v^n - B_k(u))\varphi_t \, dxdt + \int_{Q_T} (a(x, t, \nabla u^n) - a(x, t, \nabla T_k(u))\nabla\varphi \, dxdt \qquad (4.44)$$
$$+ \int_{Q_T} (H(x, t, \nabla u^n) - H(x, t, \nabla T_k(u)))\varphi \, dxdt$$
$$= \int_{Q_T} (\mu_d^n - \mu_d)\varphi \, dxdt + \int_Q (\mu_s^n - \Gamma^k)\varphi \, dxdt.$$

Using Proposition 4.1 and Proposition 4.3 we obtain from (4.44) in the sense of distribution

$$\Gamma^k = \mu_s + \omega(n,k)$$
 in  $\mathcal{D}'(Q_T)$ .

To complete the proof we have to show that the previous limit is actually tight. Let us choose without loss of generality  $\varphi \in C^1(\overline{Q_T})$  (then by density argument we show the result holds with  $\varphi \in C(\overline{Q_T})$ ). We have

$$\int_{Q_T} \Gamma^k \varphi \, dx dt = \int_{Q_T} \Gamma^k \Psi_\delta \varphi \, dx dt + \int_{Q_T} \Gamma^k (1 - \Psi_\delta) \varphi \, dx dt$$

where  $\Psi_{\delta} = \psi_{\delta}^{+} + \psi_{\delta}^{-}$  is chosen as in Lemma 4.7. Thanks to the previeous result we can write

$$\int_{Q_T} \Gamma^k \Psi_\delta \varphi \, dx dt = \int_{Q_T} \mu_s^+ \Psi_\delta \varphi \, dx dt - \int_{Q_T} \mu_s^- \Psi_\delta \varphi \, dx dt + \omega(n,k),$$

we have

$$\int_{Q_T} \mu_s^+ \Psi_\delta \varphi \, dx dt = \int_{K_\delta^+} \mu_s^+ \psi_\delta^+ \varphi \, dx dt + \int_{E^+ \setminus K_\delta^+} \mu_s^+ \psi_\delta^+ \varphi \, dx dt + \int_{Q_T} \mu_s^+ \psi_\delta^- \varphi \, dx dt,$$

since  $\psi_{\delta}^+ = 1$  on  $K_{\delta}^+$  by Lebesgue's theorem we have

$$\int_{Q_T} \mu_s^+ \Psi_\delta \varphi \, dx dt = \int_{Q_T} \mu_s^+ \varphi \, dx dt + \omega(\delta)$$

by Lemma 4.7 we obtain

$$\Big|\int_{E^+\setminus K_{\delta}^+}\mu_s^+\psi_{\delta}^+\varphi\,dxdt\Big|\leq \delta\|\varphi\|_{L^{\infty}(Q_T)},$$

and

$$\left|\int_{Q_T} \mu_s^+ \psi_{\delta}^- \varphi \, dx dt\right| \le \|\varphi\|_{L^{\infty}(Q_T)} \int_{Q_T} \psi_{\delta}^- \varphi \, d\mu_s^+ = \omega(\delta).$$

Then we otain

$$\int_{Q_T} \mu_s^+ \Psi_\delta \varphi \, dx dt = \int_{Q_T} \mu_s^+ \varphi \, dx dt + \omega(\delta).$$

Similarly we obtain

$$\int_{Q_T} \mu_s^- \Psi_\delta \varphi \, dx dt = \int_{Q_T} \mu_s^- \varphi \, dx dt + \omega(\delta) dt$$

Hence

$$\int_{Q_T} \Gamma^k \Psi_\delta \varphi \, dx dt = \int_{Q_T} \varphi \, d\mu_s + \omega(k, \delta).$$

To conclude we have to prove that

$$\int_{Q} \Gamma^{k} (1 - \Psi_{\delta}) \varphi \, dx dt = \omega(k, \delta).$$

From the definition of  $\Gamma^k$  we have

$$\int_{Q_T} (1 - \Psi_{\delta}) \varphi \, d\Gamma^k = \lim_n \left( \lim_{\sigma} \frac{1}{\sigma} \int_{\{k < |u^n| \le k + \sigma\}} sign(u^n) a(x, t, \nabla u^n) \cdot \nabla u^n (1 - \Psi_{\delta}) \varphi \right. \\ \left. + \int_{\{|u^n| \le k\}} (1 - \Psi_{\delta}) \varphi \, d\mu_s^n - \int_{\{|u^n| > k\}} (1 - \Psi_{\delta}) \varphi \, d\mu_d^n \right).$$

By Proposition 2.3 the sequence  $\mu_d^n$  is equi-diffuse, thanks to assumption (3.2) and Theorem 4.6 we deduce that

$$\left|\int_{\{|u^n|>k\}} (1-\Psi_{\delta})\varphi \,d\mu_d^n\right| \le \|\varphi\|_{L^{\infty}(Q_T)} \int_{\{|v^n|>k\gamma\}} |\mu_d^n| \,dxdt = \omega(n,k).$$

We have

$$\int_{\{|u^n| \le k\}} (1 - \Psi_{\delta}) \varphi \, d\mu_s^n = \int_{\{|u^n| \le k\}} (1 - \Psi_{\delta}) \varphi \, d\lambda_+^n - \int_{\{|u^n| \le k\}} (1 - \Psi_{\delta}) \varphi \, d\lambda_-^n,$$

and

$$\int_{\{|u^n|\leq k\}} (1-\Psi_{\delta})\varphi \,d\lambda^n_+ = \int_{\{|u^n|\leq k\}} (1-\psi^+_{\delta})\varphi \,d\lambda^n_+ - \int_{\{|u^n|\leq k\}} \psi^-_{\delta}\varphi \,d\lambda^n_+,$$

Thanks to Lemma 4.7 we obtain

$$\left|\int_{\{|u^n|\leq k\}} (1-\Psi_{\delta})\varphi \,d\lambda_+^n\right| \leq \|\varphi\|_{L^{\infty}(Q_T)} \Big(\int_{Q_T} (1-\psi_{\delta}^+)d\lambda_+^n + \int_{Q_T} \psi_{\delta}^- d\lambda_+^n\Big) = \omega(n,\delta).$$

Similarly we obtain

$$\int_{\{|u^n|\leq k\}} (1-\Psi_{\delta})\varphi \, d\lambda_{-}^n = \omega(n,\delta),$$

and then

$$\int_{\{|u^n|\leq k\}} (1-\Psi_{\delta})\varphi \, d\mu_s^n = \omega(n,\delta).$$

It remains to prove that

$$\int_{\{k < |u^n| \le k + \sigma\}} \frac{1}{\sigma} sign(u^n) a(x, t, \nabla u^n) \cdot \nabla u^n (1 - \Psi_\delta) \varphi \, dx dt = \omega(\sigma, n, k, \delta).$$

we use  $h_{k,\sigma}(u^n)(1-\Psi_{\delta})$  as test function in (4.3) we obtain

$$\int_{Q_T} B^{\star}_{h_{k,\sigma}}(u^n)(\Psi_{\delta})_t + \int_{\Omega} B^{\star}_{h_{k,\sigma}}(u^n)(T) - \int_{\Omega} B^{\star}_{h_{k,\sigma}}(u^n_0)(1 - \Psi_{\delta}(0)) \qquad (4.45)$$

$$+ \frac{1}{\sigma} \int_{\{k < |u^n| \le k + \sigma\}} a(x, t, \nabla u^n) \cdot \nabla u^n (1 - \Psi_{\delta}) - \int_{Q_T} a(x, t, \nabla u^n) \cdot \nabla \Psi_{\delta} h_{k,\sigma}(u^n)$$

$$+ \int_{Q_T} H(x, t, \nabla u^n) h_{k,\sigma}(u^n)(1 - \Psi_{\delta})$$

$$= \int_{Q_T} h_{k,\sigma}(u^n)(1 - \Psi_{\delta}) \mu^n_d + \int_{Q_T} h_{k,\sigma}(u^n)(1 - \Psi_{\delta}) \mu^n_s.$$

Using assumption (3.2), the convergence in  $L^1(Q_T)$  of  $u^n$ ,  $a(x, t, \nabla u^n)$ ,  $H(x, t, \nabla u^n)$ and the regularity of  $\Psi_{\delta}$  we obtain

$$\int_{\Omega} B_{h_{k,\sigma}}^{\star}(u^{n})(T) = \omega(\sigma, n, k),$$
$$\int_{Q_{T}} B_{h_{k,\sigma}}^{\star}(u^{n})(\Psi_{\delta})_{t} = \omega(\sigma, n, k),$$
$$\int_{\Omega} B_{h_{k,\sigma}}^{\star}(u_{0}^{n})(1 - \Psi_{\delta}(0)) = \omega(\sigma, n, k),$$
$$\int_{Q_{T}} a(x, t, \nabla u^{n}) \cdot \nabla \Psi_{\delta} h_{k,\sigma}(u^{n}) = \omega(\sigma, n, k),$$
$$\int_{Q_{T}} H(x, t, \nabla u^{n}) h_{k,\sigma}(u^{n})(1 - \Psi_{\delta}) = \omega(\sigma, n, k).$$

Thanks to Theorem 4.6 and equi-diffuse property of  $\mu_d^n$ 

$$\int_{Q_T} h_{k,\sigma}(u^n)(1-\Psi_\delta)\mu_d^n = \omega(\sigma, n, k),$$

finally by Lemma 4.7 we have

$$\int_{Q_T} h_{k,\sigma}(u^n)(1-\Psi_\delta)\mu_s^n = \omega(\sigma, n, \delta).$$

Hence we conclude that u is a renormalized solution of Problem (1.1)-(1.3).  $\Box$ 

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