# Constant Angle Spacelike Surface in de Sitter Space $S_{1}^{3}$ 

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#### Abstract

In this paper; using the angle between unit normal vector field of surfaces and a fixed spacelike axis in $R_{1}^{4}$, we develop two class of spacelike surface which are called constant timelike angle surfaces with timelike and spacelike axis in de Sitter space $S_{1}^{3}$. Moreover we give constant timelike angle tangent surfaces which are examples constant angle surfaces in de Sitter space $S_{1}^{3}$.


Key Words: Constant angle surfaces, de Sitter space, Helix.

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## 1. Introduction And Results

A constant angle curve whose tangents make constant angle with a fixed direction in ambient space is called a helix. A surface whose tangent planes makes a constant angle with a fixed vector field of ambient space is called constant angle surface. Constant angle surfaces have been studied for arbitrary dimension in Euclidean space $\mathbb{E}^{n}[13,14]$ and recently in product spaces $\mathbb{S}^{2} \times \mathbb{R}[15], \mathbb{H}^{2} \times \mathbb{R}[16]$ or different ambient spaces $\mathrm{Nil}_{3}$ [17]. In [1], Lopez and Munteanu studied constant hyperbolic angle surfaces whose unit normal timelike vector field makes a constant hyperbolic angle with a fixed timelike axis in Minkowski space $\mathbb{R}_{1}^{4}$. In the literature constant timelike and spacelike angle surface have not been investigated both in hyperbolic space $H^{3}$ and de sitter space $S_{1}^{3}$. A constant timelike and a spacelike angle surface in Hyperbolic space $H^{3}$ are developed in our paper [19]. In this paper we introduce constant timelike angle spacelike surfaces in de Sitter space $S_{1}^{3}$.

Let $x: M \longrightarrow \mathbb{R}_{1}^{4}$ be an immersion of a surface $M$ into $\mathbb{R}_{1}^{4}$. We say that $x$ is timelike (resp. spacelike, lightlike) if the induced metric on $M$ via $x$ is Lorentzian (resp. Riemannian, degenerated). If $\langle x, x\rangle=1$, then $x$ is an immersion of $S_{1}^{3}$.

Let $x: M \longrightarrow S_{1}^{3}$ be a immersion and let $\xi$ be a timelike unit normal vector field to $M$. If there exists spacelike direction $W$ such that timelike angle $\theta(\xi, U)$

[^0]is constant on $M$, then $M$ is called constant timelike angle surfaces with spacelike axis.

Let $x: M \longrightarrow S_{1}^{3}$ be a immersion and let $\xi$ be a timelike unit normal vector field to $M$. If there exists timelike direction $W$ such that timelike angle $\theta(\xi, U)$ is constant on $M$, then $M$ is called constant timelike angle surfaces with timelike axis.

## 2. Differantial Geometry of de Sitter Space $S_{1}^{3}$

In this section, Differential geometry of curves and surfaces are summarized in de Sitter space $S_{1}^{3}$. Let $\mathbb{R}_{1}^{4}$ be 4 -dimensional vector space equipped with the scalar product $\langle$,$\rangle which is defined by$

$$
\langle x, y\rangle=-x_{1} y_{1}+x_{2} y_{2}+x_{3} y_{3}+x_{4} y_{4}
$$

From now on, the constant angle surface is proposed in Minkowskian ambient space $\mathbb{R}_{1}^{4} . \mathbb{R}_{1}^{4}$ is 4 -dimensional vector space equipped with the scalar product $\langle$,$\rangle , than$ $\mathbb{R}_{1}^{4}$ is called Lorentzian 4 - space or 4 -dimensional Minkowski space. The Lorentzian norm (length) of $x$ is defined to be

$$
\|x\|=|\langle x, x\rangle|^{\frac{1}{2}} .
$$

If $\left(x_{0}^{i}, x_{1}^{i}, x_{2}^{i}, x_{3}^{i}\right)$ is the coordinate of $x_{i}$ with respect to canonical basis $\left\{e_{0}, e_{1}, e_{2}\right.$, $\left.e_{3}\right\}$ of $\mathbb{R}_{1}^{4}$, then the lorentzian cross product $x_{1} \times x_{2} \times x_{3}$ is defined by the symbolic determinant

$$
x_{1} \times x_{2} \times x_{3}=\left|\begin{array}{cccc}
-e_{0} & e_{1} & e_{2} & e_{3} \\
x_{0}^{1} & x_{1}^{1} & x_{2}^{1} & x_{3}^{1} \\
x_{0}^{2} & x_{1}^{2} & x_{2}^{2} & x_{3}^{2} \\
x_{0}^{3} & x_{1}^{3} & x_{2}^{3} & x_{3}^{3}
\end{array}\right| .
$$

One can easly see that

$$
\left\langle x_{1} \times x_{2} \times x_{3}, x_{4}\right\rangle=\operatorname{det}\left(x_{1}, x_{2}, x_{3}, x_{4}\right) .
$$

In [2], [3] and [5] Izimuya at all introduced and investigated differantial geometry of curves and surfaces Hyperbolic 3-space. If $\langle x, x\rangle>0,\langle x, x\rangle=0$ or $\langle x, x\rangle<0$ for any non-zero $x \in \mathbb{R}_{1}^{4}$, then we call that $x$ is spacelike, ligtlike or timelike ,respectively. In the rest of this section, we give background of context in [20].

Given a vector $v \in \mathbb{R}_{1}^{4}$ and a real number $c$, the hyperplane with pseudo normal $v$ is defined by

$$
H P(v, c)=\left\{x \in \mathbb{R}_{1}^{4} \mid\langle x, v\rangle=c\right\}
$$

We say that $H P(v, c)$ is a spacelike hyperplane, timelike hyperplane or lightlike hyperplane if $v$ is timelike, spacelike or lightlike respectively in [20]. We have following three types of pseudo-spheres in $\mathbb{R}_{1}^{4}$ :

$$
\begin{aligned}
& \text { Hyperbolic-3 space : } H^{3}(-1)=\left\{x \in \mathbb{R}_{1}^{4} \mid\langle x, x\rangle=-1, x_{0} \geq 1\right\}, \\
& \text { de Sitter 3- space: } S_{1}^{3}=\left\{x \in \mathbb{R}_{1}^{4} \mid\langle x, x\rangle=1\right\}, \\
& \quad \text { (open) lightcone: } L C^{*}=\left\{x \in \mathbb{R}_{1}^{4} /\{0\} \mid\langle x, x\rangle=0, x_{0}>0\right\} .
\end{aligned}
$$

We also define the lightcone 3 -sphere

$$
S_{+}^{3}=\left\{x=\left(x_{0}, x_{1}, x_{2}, x_{3}\right) \mid\langle x, x\rangle=0, x_{0}=1\right\}
$$

A hypersurface given by the intersection of $S_{1}^{3}$ with a spacelike (resp.timelike) hyperplane is called an elliptic hyperquadric (resp. hyperbolic hyperquadric). If $c \neq 0$ and $H P(v, c)$ is lightlike, then $H P(v, c) \cap S_{1}^{3}$ is a de Sitter horosphere, [20].

Let $U \subset \mathbb{R}^{2}$ be open subset, and let $x: U \rightarrow S_{1}^{3}$ be an embedding. If the vector subspace $\tilde{U}$ whih generated by $\left\{x_{u_{1}}, x_{u_{2}}\right\}$ is spacelike, then $x$ is called spacelike surface, if $\tilde{U}$ contain at least a timelike vector field, then $x$ is called timelike surface in $S_{1}^{3}$.

In point of view Kasedou [20], we construct the extrinsic differential geometry on curves in $S_{1}^{3}$. Since $S_{1}^{3}$ is a Riemannian manifold, the regular curve $\gamma: I \rightarrow S_{1}^{3}$ is given by arclength parameter.

Theorem 2.1. i) If $\gamma: I \rightarrow S_{1}^{3}$ is a spacelike curve with unit speed, then FrenetSerre type formulae is obtained

$$
\left\{\begin{aligned}
\gamma^{\prime}(s) & =t(s) \\
t^{\prime}(s) & =\kappa_{d}(s) n(s)-\gamma(s) \\
n^{\prime}(s) & =-\kappa_{d}(s) t(s)-\tau_{d}(s) e(s) \\
e^{\prime}(s) & =-\tau_{d}(s) n(s)
\end{aligned}\right.
$$

where $\kappa_{d}(s)=\left\|t^{\prime}(s)+\gamma(s)\right\|$ and $\tau_{d}(s)=-\frac{\operatorname{det}\left(\gamma(s), \gamma^{\prime}(s), \gamma^{\prime \prime}(s), \gamma^{\prime \prime \prime}(s)\right)}{\left(\kappa_{d}(s)\right)^{2}}$.
ii) If $\gamma: I \rightarrow S_{1}^{3}$ is a timelike curve with unit speed, then Frenet-Serre type formulae is obtained

$$
\begin{cases}\gamma^{\prime}(s) & =t(s) \\ t^{\prime}(s) & =\kappa_{d}(s) n(s)+\gamma(s) \\ n^{\prime}(s) & =-\kappa_{d}(s) t(s)+\tau_{d}(s) e(s) \\ e^{\prime}(s) & =-\tau_{d}(s) n(s)\end{cases}
$$

where $\kappa_{d}(s)=\left\|t^{\prime}(s)-\gamma(s)\right\|$ and $\tau_{d}(s)=-\frac{\operatorname{det}\left(\gamma(s), \gamma^{\prime}(s), \gamma^{\prime \prime}(s), \gamma^{\prime \prime \prime}(s)\right)}{\left(\kappa_{d}(s)\right)^{2}}$.
It is easily see that $\kappa_{d}(s)=0$ if and only if there exists a lightlike vector $c$ such that $\gamma(s)-c$ is a geodesic.

Now we give extrinsic differential geometry on surfaces in $S_{1}^{3}$ due to Kasedou [20].

Let $U \subset \mathbb{R}^{2}$ is an open subset and $x: U \rightarrow S_{1}^{3}$ is a regular surface $M=x(U)$. If $e(u)$ is defined as follows

$$
e(u)=\frac{x(u) \wedge x_{u_{1}}(u) \wedge x_{u_{2}}(u)}{\left\|x(u) \wedge x_{u_{1}}(u) \wedge x_{u_{2}}(u)\right\|}
$$

then

$$
\langle e, x\rangle \equiv\left\langle e, x_{u_{i}}\right\rangle \equiv 0,\langle e, e\rangle=-1
$$

where $x_{u_{i}}=\frac{\partial x}{\partial u_{i}}$. Thus there is de Sitter Gauss image of $x$ which is defined by mapping

$$
E: U \rightarrow S_{1}^{3}, E(u)=e(u) .
$$

The lightcone Gauss image of $x$ is defined by map

$$
L^{ \pm}: U \rightarrow L C^{*}, L^{ \pm}(u)=x(u) \pm e(u) .
$$

Let $d x\left(u_{0}\right)$ and $1_{T_{p} M}$ be identify mapping on the tangent space $T_{p} M$. So derivate $d x\left(u_{0}\right)$ can be identified with $T_{p} M$ relate to identification of $U$ and $M$. That is

$$
d L^{ \pm}\left(u_{0}\right)=1_{T_{p} M} \pm d E\left(u_{0}\right)
$$

The linear transformation

$$
S_{p}^{ \pm}:=-d L^{ \pm}\left(u_{0}\right): T_{p} M \rightarrow T_{p} M
$$

and

$$
A_{p}:=-d E\left(u_{0}\right): T_{p} M \rightarrow T_{p} M
$$

is called the hyperbolic shape operator and de Sitter shape operator of $M$ at $p=$ $x\left(u_{o}\right)$.

Let $\bar{K}_{i}^{ \pm}(p)$ and $K_{i}(p),(i=1,2)$ be the eigenvalues of $S_{p}^{ \pm}$and $A_{p}$. Since

$$
S_{p}^{ \pm}=-1_{T_{p} M} \pm A_{p},
$$

$S_{p}^{ \pm}$and $A_{p}$ have same eigenvectors and relations

$$
\bar{K}_{i}^{ \pm}(p)=-1 \pm K_{i}(p) .
$$

$\bar{K}_{i}^{ \pm}(p)$ and $K_{i}(p),(i=1,2)$ are called hyperbolic and de Sitter principal curvatures of $M$ at $p=x\left(u_{0}\right)$.

Let $\gamma(s)=x\left(u_{1}(s), u_{2}(s)\right)$ be a unit speed curve on $M$, with $p=\gamma\left(u_{1}\left(s_{0}\right)\right.$, $\left.u_{2}\left(s_{0}\right)\right)$. We consider the hyperbolic curvature vector $k(s)=t^{\prime}(s)-\gamma(s)$ and the de Sitter normal curvature

$$
K_{n}^{ \pm}\left(s_{0}\right)=\left\langle k\left(s_{0}\right), L^{ \pm}\left(u_{1}\left(s_{0}\right), u_{2}\left(s_{0}\right)\right)\right\rangle=\left\langle t^{\prime}\left(s_{0}\right), L^{ \pm}\left(u_{1}\left(s_{0}\right), u_{2}\left(s_{0}\right)\right)\right\rangle+1
$$

of $\gamma(s)$ at $p=\gamma\left(s_{0}\right)$. The de Sitter normal curvature depends only on the point $p$ and the unit tangent vector of $M$ at $p$ analogous to the Euclidean case. Hyperbolic normal curvature of $\gamma(s)$ is defined to be

$$
\bar{K}_{n}^{ \pm}(s)=K_{n}^{ \pm}(s)-1 .
$$

The Hyperbolic Gauss curvature of $M=x(U)$ at $p=x\left(u_{0}\right)$ is defined to be

$$
K_{h}^{ \pm}\left(u_{0}\right)=\operatorname{det} S_{p}^{ \pm}=\bar{K}_{1}^{ \pm}(p) \bar{K}_{2}^{ \pm}(p) .
$$

The Hyperbolic mean curvature of $M=x(u)$ at $p=x\left(u_{0}\right)$ is defined to be

$$
H_{h}^{ \pm}\left(u_{0}\right)=\frac{1}{2} \operatorname{Trace} S_{p}^{ \pm}=\frac{\bar{K}_{1}^{ \pm}(p)+\bar{K}_{2}^{ \pm}(p)}{2}
$$

The extrinsic (de Sitter) Gauss curvature is defined to be

$$
K_{e}\left(u_{0}\right)=\operatorname{det} A_{p}=K_{1}(p) K_{2}(p),
$$

and the de Sitter mean curvature is

$$
H_{d}\left(u_{0}\right)=\frac{1}{2} \operatorname{Trace} A_{p}=\frac{K_{1}(p)+K_{2}(p)}{2} .
$$

## 3. Constant Timelike Angle Spacelike Surfaces

Let us show the space of the tangent vector fields on $M$ with $X(M)$ and denote the Levi-Civita connections of $\mathbb{R}_{1}^{4}, S_{1}^{3}$ and $M$ by $\overline{\bar{D}}, \bar{D}$ and $D$. Then for each $X, Y \in X(M)$, we have

$$
D_{X} Y=\left(\overline{\bar{D}}_{X} Y\right)^{T}, \widetilde{V}(X, Y)=\left(\overline{\bar{D}}_{X} Y\right)^{\perp}
$$

and

$$
\begin{equation*}
\overline{\bar{D}}_{X} Y=\bar{D}_{X} Y-\langle X, Y\rangle x, \overline{\bar{D}}_{X} Y=D_{X} Y+\widetilde{V}(X, Y) \tag{3.1}
\end{equation*}
$$

where the superscript ${ }^{T}$ and ${ }^{\perp}$ denote the tangent and normal component of $\overline{\bar{D}}_{X} Y$. (3.1) equation is called the Gauss formula of $S_{1}^{3}$ and $M$.

If $\xi$ is a normal vector field of $M$ on $S_{1}^{3}$, then the Weingarten Endomorphism $A_{\xi}(X)$ and $B_{x}(X)$ are denoted by the tangent components of $-\overline{\bar{D}}_{X} \xi$ and $-\overline{\bar{D}}_{X} x$. So the Weingarten equations of the vector field $\xi$ and $x$ is like

$$
\left\{\begin{array}{l}
A_{\xi}(X)=-\overline{\bar{D}}_{X} \xi-\left\langle\overline{\bar{D}}_{X} x, \xi\right\rangle x  \tag{3.2}\\
B_{x}(X)=-\overline{\bar{D}}_{X} x-\left\langle\overline{\bar{D}}_{X} x, \xi\right\rangle \xi
\end{array}\right.
$$

It is obvious that $A_{\xi}(X)$ and $B_{x}(X)$ are linear and self adjoint map for each $p \in M$. That is

$$
\left\langle A_{\xi}(X), Y\right\rangle=\left\langle X, A_{\xi}(Y)\right\rangle \text { and }\left\langle B_{x}(X), Y\right\rangle=\left\langle X, B_{x}(Y)\right\rangle
$$

The eigenvalues $K_{i}(p)$ and $\tilde{K}_{i}(p)$ of $\left(A_{\xi}\right)_{p}$ are called the principal curvature of $M$ on $S_{1}^{3}$. The eigenvalues $\tilde{K}_{i}(p)$ of $\left(B_{x}\right)_{p}$ are called the principal curvature of $M$ in $\mathbb{R}_{1}^{4}$. Also, for $X, Y \in X(M)$ we have

$$
\left\langle A_{\xi}(X), Y\right\rangle=\langle\widetilde{V}(X, Y), \xi\rangle \quad, \quad\left\langle B_{x}(X), Y\right\rangle=\langle\widetilde{V}(X, Y), x\rangle
$$

Since $\widetilde{V}(X, Y)$ is second fundamental form of $M$ on $\mathbb{R}_{1}^{4}$, so we can write as follows

$$
\widetilde{V}(X, Y)=-\langle\widetilde{V}(X, Y), \xi\rangle \xi+\langle\widetilde{V}(X, Y), x\rangle x
$$

and

$$
\tilde{V}(X, Y)=-\left\langle A_{\xi}(X), Y\right\rangle \xi+\left\langle B_{x}(X), Y\right\rangle x
$$

Let $\left\{v_{1}, v_{2}\right\}$ be a base of $T_{p} M$ tangent plane and let us denote

$$
\begin{align*}
a_{i j} & =\left\langle\widetilde{V}\left(v_{i}, v_{j}\right), \xi\right\rangle=\left\langle A_{\xi}\left(v_{i}\right), v_{j}\right\rangle  \tag{3.3}\\
b_{i j} & =\left\langle\widetilde{V}\left(v_{i}, v_{j}\right), x\right\rangle=\left\langle B_{x}\left(v_{i}\right), v_{j}\right\rangle \tag{3.4}
\end{align*}
$$

Therefore

$$
\overline{\bar{D}}_{X} Y=D_{X} Y+\widetilde{V}(X, Y)
$$

and also since

$$
\bar{D}_{X} Y=D_{X} Y-\left\langle A_{\xi}(X), Y\right\rangle \xi \text { and } \bar{D}_{X} Y=\bar{D}_{X} Y-\langle X, Y\rangle x
$$

we obtain

$$
\overline{\bar{D}}_{X} Y=D_{X} Y-\left\langle A_{\xi}(X), Y\right\rangle \xi-\langle X, Y\rangle x
$$

On the other hand for $\left\{v_{1}, v_{2}\right\}$ base, we get

$$
\begin{equation*}
\overline{\bar{D}}_{v_{i}} v_{j}=D_{v_{i}} v_{j}-a_{i j} \xi-\left\langle v_{i}, v_{j}\right\rangle x \tag{3.5}
\end{equation*}
$$

If this basis is orthonormal, then we have from (3.1) and (3.2)

$$
\begin{align*}
& \overline{\bar{D}}_{v_{i}} v_{j}=D_{v_{i}} v_{j}-a_{i j} \xi,  \tag{3.6}\\
& \overline{\bar{D}}_{v_{i}} \xi=-a_{i 1} v_{1}-a_{i_{2}} v_{2},  \tag{3.7}\\
& \overline{\bar{D}}_{v_{i}} x=-b_{i 1} v_{1}-b_{i 2} v_{2} . \tag{3.8}
\end{align*}
$$

### 3.1. Constant Timelike Angle Surfaces With Spacelike Axis

Definition 3.1. Let $U \subset \mathbb{R}^{2}$ be open set,let $x: U \rightarrow S_{1}^{3}$ be an embedding where $M=x(U)$. Let $x: M \rightarrow S_{1}^{3}$ and $\xi$ is timelike unit normal vector field on $M$, if there exist a constant spacelike vector $W$ which has a constant timelike angle with $\xi$, then $M$ is called constant timelike angle surface with spacelike axis.

Since our surface is a spacelike surface, $\left\{x_{u}, x_{v}\right\}$ tangent vectors must be spacelike vectors. Let $M$ be a spacelike surface with constant angle with spacelike axis and $\xi$ is unit normal vector of $M$ on $S_{1}^{3}$. Let us denote that the timelike angle between timelike vector $\xi$ and spacelike vector $W$ with $\theta$. That is from [11]

$$
\langle\xi, W\rangle=\sin h(-\theta) .
$$

If timelike angle $\theta=0$, then $\xi=W$. Throughout this section, without loss of generality we assume that $\theta \neq 0$. If $W^{T}$ is the projection of $W$ on the tangent plane of $M$, then we decompose $W$ as

$$
W=W^{T}+W^{N}
$$

So that we write

$$
W=W^{T}+\lambda_{1} \xi+\lambda_{2} x .
$$

If we take inner product of both sides of this inequality first with $\xi$, then with $x$

$$
\lambda_{1}=-\sin h(-\theta), \lambda_{2}=\langle W, x\rangle
$$

On the other hand, since $W$ and $x$ are two spacelike vector fields, then we can use define of the spacelike and timelike angle between $W$ and $x$.

Theorem 3.2. i) If $\varphi$ is the spacelike angle between spacelike vectors $W, x$ then we can write for [11]

$$
W=\sqrt{\sinh ^{2} \theta+\sin ^{2} \varphi} e_{1}+(\sinh \theta) \xi+\cos \varphi x
$$

and de Sitter projection $W_{d}$ of $W$ as follows

$$
\begin{equation*}
W_{d}=\sqrt{\sinh ^{2} \theta+\sin ^{2} \varphi} e_{1}+(\sinh \theta) \xi \tag{3.9}
\end{equation*}
$$

ii)If $\varphi$ is timelike angle between spacelike vectors $W$ and $x$, then we can write

$$
W=\sqrt{\left|\cosh ^{2} \theta-\cosh ^{2} \varphi\right|} e_{1}+(\sinh \theta) \xi-(\cosh \varphi) x
$$

and de sitter projection $W_{d}$ of $W$ as follows

$$
\begin{equation*}
W_{d}=\sqrt{\left|\cosh ^{2} \theta-\cosh ^{2} \varphi\right|} e_{1}+(\sinh \theta) \xi \tag{3.10}
\end{equation*}
$$

Let $e_{1}=\frac{W^{T}}{\left\|W^{T}\right\|}$ and let consider $e_{2}$ be a unit vector field on $M$ orthogonal to $e_{1}$. Then we have an oriented orthonormal basis $\left\{e_{1}, e_{2}, \xi, x\right\}$ for $\mathbb{R}_{1}^{4}$. Since $W_{d}$ is constant vector field on $S_{1}^{3}$ and $\overline{\bar{D}}_{e_{2}} W_{d}=\bar{D}_{e_{2}} W_{d}=0$, we have

$$
\begin{equation*}
\sqrt{\sinh ^{2} \theta+\sin ^{2} \varphi} \overline{\bar{D}}_{e_{2} e_{1}+(\sinh \theta)} \overline{\bar{D}}_{e_{2}} \xi=0 \tag{3.11}
\end{equation*}
$$

By (3.11), we obtain

$$
\sqrt{\sinh ^{2} \theta+\sin ^{2} \varphi}\left\langle\overline{\bar{D}}_{e_{2}} e_{1}, \xi\right\rangle+(\sinh \theta)\left\langle\overline{\bar{D}}_{e_{2}} \xi, \xi\right\rangle=0
$$

or

$$
\sqrt{\sinh ^{2} \theta+\sin ^{2} \varphi} a_{21}=0
$$

Since $\sqrt{\sinh ^{2} \theta+\sin ^{2} \varphi} \neq 0$, we conclude $a_{21}=a_{12}=0$. Using (3.7) in (3.11), we get

$$
\begin{equation*}
\overline{\bar{D}}_{e_{2}} e_{1}=\frac{\sinh \theta}{\sqrt{\sinh ^{2} \theta+\sin ^{2} \varphi}} a_{22} e_{2} . \tag{3.12}
\end{equation*}
$$

Similarly, since $W_{d}$ is a constant vector field on $S_{1}^{3}$, then we have

$$
\begin{equation*}
\bar{D}_{e_{1}} W_{d}=0 \text { and } \overline{\bar{D}}_{e_{1}} W_{d}=-\sqrt{\sinh ^{2} \theta+\sin ^{2} \varphi x} \tag{3.13}
\end{equation*}
$$

By (3.9), we obtain

$$
\begin{equation*}
\overline{\bar{D}}_{e_{1}} W_{d}=\sqrt{\sinh ^{2} \theta+\sin ^{2} \varphi} \overline{\bar{D}}_{e_{1}} e_{1}+\sinh \theta \overline{\bar{D}} e_{1} \xi \tag{3.14}
\end{equation*}
$$

By (3.13) and (3.14), we conclude that

$$
\begin{equation*}
\sqrt{\sinh ^{2} \theta+\sin ^{2} \varphi} \overline{\bar{D}}_{e_{1}} e_{1}+\sinh \theta \overline{\bar{D}} e_{1} \xi=-\sqrt{\sinh ^{2} \theta+\sin ^{2} \varphi} x \tag{3.15}
\end{equation*}
$$

By (3.15), we get

$$
\sqrt{\sinh ^{2} \theta+\sin ^{2} \varphi}\left\langle\overline{\bar{D}}_{e_{1}} e_{1}, \xi\right\rangle=0
$$

or

$$
\sqrt{\sinh ^{2} \theta+\sin ^{2} \varphi} a_{11}=0
$$

Since $\sqrt{\sinh ^{2} \theta+\sin ^{2} \varphi} \neq 0$, we conclude $a_{11}=0$. Also , using (3.7) in (3.15), we obtain

$$
\begin{equation*}
\overline{\bar{D}}_{e_{1} e_{1}}=-x \tag{3.16}
\end{equation*}
$$

Now we have proved the following theorem.
Theorem 3.3. If $D$ is Levi-Civita connection for a constant timelike angle with spacelike axis spacelike surface in $S_{1}^{3}$ is given by

$$
\begin{array}{ll}
D_{e_{1}} e_{1}=0, & D_{e_{2}} e_{1}=\frac{\sinh \theta}{\sqrt{\sinh ^{2} \theta+\sin ^{2} \varphi}} a_{22} e_{2} \\
D_{e_{1}} e_{2}=0, & D_{e_{2}} e_{2}=\frac{-\sinh \theta}{\sqrt{\sinh ^{2} \theta+\sin ^{2} \varphi}} a_{22} e_{1}
\end{array}
$$

Corollary 3.4. Let $M$ be a spacelike surface which is a constant timelike angle with spacelike axis on $S_{1}^{3}$. Then, there exist local coordinates $u$ and $v$ such that the metric on $M$ writes as $\langle\rangle=,d u^{2}+\beta^{2} d v^{2}$, where $\beta=\beta(u, v)$ is a smooth function on $M$, i.e. the coefficients of the first fundamental form are $E=1, F=0, G=\beta^{2}$.

Now we find the $x=x(u, v)$ parametrization of the surface $M$ with respect to the metric $\langle\rangle=,d u^{2}+\beta^{2} d v^{2}$ on $M$. By the above parametrization $x(u, v)$ can obtain the following corollary.
Corollary 3.5. There exist an equation system for constant timelike angle with spacelike axis spacelike surface on $S_{1}^{3}$ which is

$$
\left\{\begin{align*}
x_{\mathrm{u} u} & =-x  \tag{3.17}\\
x_{u v} & =\frac{\beta_{u}}{\beta} x_{v} \\
x_{v v} & =-\beta \beta_{u} x_{u}+\frac{\beta_{v}}{\beta} x_{v}-\beta^{2} a_{22} \xi-\beta^{2} x
\end{align*}\right.
$$

Corollary 3.6. Let $\xi$ be unit normal vector of the constant timelike angle with spacelike axis spacelike surface $M$. Then the equation below hold

$$
\left\{\begin{array}{l}
\xi_{u}=\overline{\bar{D}}_{x_{u}} \xi=0  \tag{3.18}\\
\xi_{v}=\overline{\bar{D}}_{x_{v}} \xi=-a_{22} x_{v} .
\end{array}\right.
$$

Since $\xi_{u v}=\xi_{v u}=0$, we have $\overline{\bar{D}}_{x_{u}}\left(-a_{22} x_{v}\right)=0$. Using $a_{12}=0, \overline{\bar{D}}_{x_{u}} x_{v}=$ $\overline{\bar{D}}_{x_{v}} x_{u}$ and Theorem 2.1, we obtain

$$
\begin{equation*}
\left(a_{22}\right)_{u}+\frac{\sinh \theta}{\sqrt{\sinh ^{2} \theta+\sin ^{2} \varphi}}\left(a_{22}\right)^{2}=0 \tag{3.19}
\end{equation*}
$$

So that

$$
\begin{equation*}
\left(a_{22}\right)_{u}+\frac{\beta_{u}}{\beta} a_{22}=0 \tag{3.20}
\end{equation*}
$$

and than we get obtain

$$
\begin{equation*}
\left(\beta a_{22}\right)_{u}=0 . \tag{3.21}
\end{equation*}
$$

By (3.21), we see that there exist a smooth function $\psi=\psi(v)$ depending on $v$ such that

$$
\begin{equation*}
\beta a_{22}=\psi(v) \tag{3.22}
\end{equation*}
$$

Proposition 3.7. Let $x=x(u, v)$ be parametrization of a spacelike surface which is constant timelike angle with spacelike axis on $S_{1}^{3}$. If $a_{22}=0$ on $M$, then the $x$ describes an flat plane of de Sitter space $S_{1}^{3}$.

Proof. If $a_{22}=0$ on $M$, then by (3.18)

$$
\left\{\begin{array}{l}
\xi_{u}=0 \\
\xi_{v}=0
\end{array}\right.
$$

This imply we have $\xi$ is a constant vector field which normal vector is $M$ surface. Thus $x=x(u, v)$ is de-Sitter plane in $S_{1}^{3}$.

From now on, we are going to assume that $a_{22} \neq 0$. By solving equation (3.19), we obtain a function $\alpha=\alpha(v)$ such that

$$
a_{22}=\frac{\sqrt{\sinh ^{2} \theta+\sin ^{2} \varphi}}{u \sinh \theta+\alpha(v)} .
$$

Therefore by (3.22) , we obtain

$$
\beta(u, v)=\frac{\psi(v)}{\sqrt{\sinh ^{2} \theta+\sin ^{2} \varphi}}(u \sinh \theta+\alpha(v))
$$

Consequently,

$$
\begin{aligned}
& x_{\mathrm{u} u}=-x, \\
& x_{u v}=\frac{\sinh \theta \psi(v)}{u \sinh \theta+\alpha(v)} x_{v}, \\
& x_{v v}=\left[\frac{-\psi^{2}(v) \sinh \theta(u \sinh \theta+\alpha(v))}{\sinh ^{2} \theta+\sin ^{2} \varphi}\right] x_{u}+\left[\frac{\psi^{\prime}(v)}{\psi(v)}+\frac{\alpha^{\prime}(v)}{u \sinh \theta+\alpha(v)}\right] x_{v} \\
& -\left[\frac{\psi^{2}(v)(u \sinh \theta+\alpha(v))}{\left.\sqrt{\sinh ^{2} \theta+\sin ^{2} \varphi}\right] \xi-\left[\frac{\psi^{2}(v)(u \sinh \theta+\alpha(v))^{2}}{\sinh ^{2} \theta+\sin ^{2} \varphi}\right] x .} .\right.
\end{aligned}
$$

Here, if we spesificly choose $\psi(v)=e^{v} \sqrt{\sinh ^{2} \theta+\sin ^{2} \varphi}$ and $\alpha(v)=e^{-v}$, then this equation system becomes

$$
\left\{\begin{align*}
x_{\mathrm{u} u} & =-x  \tag{3.23}\\
x_{u v} & =\frac{e^{v} \sinh \theta}{1+u e^{v} \sinh \theta} x_{v} \\
x_{v v} & =-e^{v} \sinh \theta\left(u e^{v} \sinh \theta+1\right) x_{u}+\frac{u e^{v} \sinh \theta}{u e^{v} \sinh \theta+1} x_{v}- \\
& -e^{v} \sqrt{\sinh ^{2} \theta+\sin ^{2} \varphi}\left(u e^{v} \sinh \theta+1\right) \xi-\left(u e^{v} \sinh \theta+1\right)^{2} x
\end{align*}\right.
$$

Now we have the following Theorem.
Theorem 3.8. If $M$ is satisfying (3.23), then there exist local coordinates $u$ and $v$ on $M$ with having the parametrization

$$
\begin{equation*}
x_{i}(u, v)=\left(\frac{-c_{1 i}(v)}{2 e^{v} \sinh \theta\left(u e^{v} \sinh \theta+1\right)^{2}}+c_{2 i}(v)\right) \quad, \quad i=1,2,3,4 \tag{3.24}
\end{equation*}
$$

Proof. From (3.23), the proof is clear.
Example 3.9. We can calculate Gauss and mean curvature of a spacelike surface with constant angle spacelike axis in de Sitter space $S_{1}^{3}$. Since

$$
\overline{\bar{D}}_{X} \xi=\bar{D}_{X} \xi
$$

we can write

$$
\overline{\bar{D}}_{v_{i}} \xi=\left\langle\overline{\bar{D}}_{v_{i}} \xi, v_{1}\right\rangle v_{1}+\left\langle\overline{\bar{D}}_{v_{i}} \xi, v_{2}\right\rangle v_{2} .
$$

Thus from (3.7), we have

$$
\overline{\bar{D}}_{v_{i}} \xi=-a_{i 1} v_{1}-a_{i_{2}} v_{2}
$$

From $A_{\xi}\left(v_{i}\right)=-\overline{\bar{D}}_{v_{i}} \xi$ and $a_{21}=a_{12}=0, a_{11}=0$, we obtain

$$
A_{\xi}=\left(\begin{array}{cc}
0 & 0 \\
0 & a_{22}
\end{array}\right)
$$

Since eigenvalues of linear transformation $A_{p}: T_{p} M \rightarrow T_{p} M$ are principal curvatures of $M$ at $p$, we obtain the following principal curvatures of $M$

$$
K_{1}(p)=0 \text { and } K_{2}(p)=a_{22}
$$

Hence Gauss and mean curvature of $M$ at $p$ are

$$
\begin{gathered}
K_{e}(p)=0 \\
H_{d}(p)=\frac{1}{2} a_{22},
\end{gathered}
$$

where $a_{22}$ is

$$
a_{22}=\frac{\sqrt{\sinh ^{2} \theta+\sin ^{2} \varphi}}{e^{-v}+u \sinh \theta}
$$

Remark 3.10. If we consider

$$
W_{d}=\sqrt{\cosh ^{2} \theta-\cosh ^{2} \varphi} e_{1}+\sinh \theta \xi
$$

the constant direction of spacelike surface with constant timelike angle in de Sitter space $S_{1}^{3}$, then we obtain similar results in Theorem 3.2-i.

Remark 3.11. If the $W_{d}$ constant direction of spacelike surface is chosen a timelike vector, then we obtain similar resultys in chapter 3.1

## 4. Constant Timelike Angle Tangent Surfaces

### 4.1. Tangent Surface with Spacelike Axis

In this section we will focus on constant timelike angle spacelike tangent surfaces with spacelike axis in de Sitter space $S_{1}^{3}$. ( see [2] and [6] for the Minkowski ambient space and Euclidean ambient space,respectively). Let $\alpha: I \rightarrow S_{1}^{3} \subset \mathbb{R}_{1}^{4}$ be a regular spacelike curve given by arc-length. We define the tangent surface $M$, which is generated by $\alpha$, with

$$
\begin{equation*}
x(s, t)=\alpha(s) \cos t+\alpha^{\prime}(s) \sin t, \quad(s, t) \in I \times \mathbb{R} \tag{4.1}
\end{equation*}
$$

The tangent plane at a point $(s, t)$ of $M$ is spanned by $\left\{x_{s}, x_{t}\right\}$, where

$$
\begin{cases}x_{s} & =\alpha^{\prime}(s) \cos t+\alpha^{\prime \prime}(s) \sin t  \tag{4.2}\\ x_{t} & =-\alpha(s) \sin t+\alpha^{\prime}(s) \cos t\end{cases}
$$

By computing the coefficients of first fundamental form $\{E, F, G\}$ of $M$ with respect to basis $\left\{x_{s}, x_{t}\right\}$, we get

$$
\left\{\begin{array}{l}
E=\left\langle x_{s}, x_{s}\right\rangle=1+\kappa_{d}^{2}(s) \sin ^{2} t, \\
F=\left\langle x_{s}, x_{t}\right\rangle=1 \\
G=\left\langle x_{t}, x_{t}\right\rangle=1 .
\end{array}\right.
$$

Hence we have

$$
E G-F^{2}=\kappa_{d}^{2}(s) \sin ^{2} t
$$

Then, since $E G-F^{2}>0$, it is obvious that $M$ is a spacelike surface. From Frenet-Serre type formulae, we obtain

$$
\begin{cases}x(s, t) & =\alpha(s) \cos t+t(s) \sin t  \tag{4.3}\\ x_{s}(s, t) & =\alpha(s) \sin t+t(s) \cos t+n(s) \kappa_{d}(s) \sin t \\ x_{t}(s, t) & =-\alpha(s) \sin t+t(s) \cos t\end{cases}
$$

Now let us calculate normal vector of $M$. As we already know the normal vector of $M$ is

$$
\begin{equation*}
e=\frac{x \wedge x_{s} \wedge x_{t}}{\left\|x \wedge x_{s} \wedge x_{t}\right\|} \tag{4.4}
\end{equation*}
$$

Then, since

$$
x \wedge x_{s} \wedge x_{t}=-\left(\alpha \wedge \alpha^{\prime} \wedge \alpha^{\prime \prime}\right) \sin t
$$

and

$$
\left\|x \wedge x_{s} \wedge x_{t}\right\|=\left|\kappa_{d} \sin t\right| \quad, \quad \kappa_{d} \neq 0
$$

we find

$$
\begin{equation*}
e= \pm \frac{\alpha \wedge \alpha^{\prime} \wedge \alpha^{\prime \prime}}{\left|\kappa_{d}\right|} \tag{4.5}
\end{equation*}
$$

Let us find $W_{d}$ direction of constant timelike angle with spacelike axis surface $M$. Since (3.9) and

$$
e_{1}=\frac{x_{s}}{\left\|x_{s}\right\|} \text { and }\left\|x_{s}\right\|=\sqrt{1+\kappa_{d}^{2} \sin ^{2} t} .
$$

we get

$$
\begin{align*}
W_{d} & =\sin t \sqrt{\frac{\sinh ^{2} \theta-\sinh ^{2} \varphi}{1+\kappa_{d}^{2} \sin ^{2} t}} \alpha(s)+\cos t \sqrt{\frac{\sinh ^{2} \theta-\sinh ^{2} \varphi}{1+\kappa_{d}^{2} \sin ^{2} t}} t(s)+  \tag{4.6}\\
& +\kappa_{d}(s) \sin t \sqrt{\frac{\sinh ^{2} \theta-\sinh ^{2} \varphi}{1+\kappa_{d}^{2} \sin ^{2} t}} n(s)+e(s) \cosh \theta
\end{align*}
$$

Theorem 4.1. Let $\alpha: I \rightarrow S_{1}^{3} \subset \mathbb{R}_{1}^{4}$ be a curve with $\kappa_{d} \neq 0$. If $x(s, t)$ tangent surface is constant timelike angle surface with spacelike axis, then $\alpha$ curve is planarly.
Proof. Suppose that $x(s, t)$ tangent surface is constant timelike angle surface with spacelike axis such that $\alpha$ is a curve with $\kappa_{d} \neq 0$. Since

$$
\xi=\frac{x \wedge x_{s} \wedge x_{t}}{\left\|x \wedge x_{s} \wedge x_{t}\right\|}=e
$$

there exsist a $\theta>0$ real number such that

$$
\left\langle\xi, W_{d}\right\rangle=\left\langle e(s), W_{d}\right\rangle=\cosh \theta
$$

If we differentiate the both sides of the last equation with respect to $s$ then we get that

$$
\left\langle e^{\prime}(s), W_{d}\right\rangle=0
$$

By the way we know that from Frenet-Serret equation system

$$
e^{\prime}(s)=\tau_{d}(s) n(s)
$$

Hence we get

$$
\begin{equation*}
\left\langle n(s), W_{d}\right\rangle=0 \text { or } \tau_{d}(s)=0 . \tag{4.7}
\end{equation*}
$$

If in equation (4.7) $\left\langle n(s), W_{d}\right\rangle=0$ then scalar producting of (4.6) equation with $n(s)$ that we have $t=0$. This is contradict with definition of tangent surface. Therefore using equation (4.7) $\tau_{d}(s)=0$ is obvious. It means that $\alpha$ is planarly line.

Example 4.2. Let $\alpha: I \rightarrow S_{1}^{3} \subset \mathbb{R}_{1}^{4}$ be a regular curve given by arc-length

$$
\alpha(s)=\left(s \sinh (\arccos h s), s \cosh (\arccos h s), \sqrt{1-s^{2}}, 0\right) .
$$

Since the tangent surface $M$ generated by $\alpha$ as the surface parametrized by

$$
x(s, t)=\alpha(s) \cos t+\alpha^{\prime}(s) \sin t, \quad(s, t) \in I \times \mathbb{R}
$$

The picture of the Stereographic projection of tangent surface appear in Figure 1


Figure 1:

Remark 4.3. If we consider

$$
W_{d}=\sqrt{\left|\cosh ^{2} \theta-\cosh ^{2} \varphi\right|} e_{1}+\sinh \theta \xi
$$

then we will get similar result.
Remark 4.4. If the $W_{d}$ constant direction of spacelike surface is chosen timelike, then we obtain similar results in chapter 4.1

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