



## $\mathcal{D}$ –Tangent Surfaces of Timelike Biharmonic $\mathcal{D}$ –Helices according to Darboux Frame on Non-degenerate Timelike Surfaces in the Lorentzian Heisenberg Group $\mathbb{H}$

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ABSTRACT: In this paper, we study  $\mathcal{D}$ –tangent surfaces of timelike biharmonic  $\mathcal{D}$ –helices according to Darboux frame on non-degenerate timelike surfaces in the Lorentzian Heisenberg group  $\mathbb{H}$ . We obtain parametric equation  $\mathcal{D}$ –tangent surfaces of timelike biharmonic  $\mathcal{D}$ –helices in the Lorentzian Heisenberg group  $\mathbb{H}$ . Moreover, we illustrate the figure of our main theorem.

Key Words:  $\mathcal{D}$ –Biharmonic curve,  $\mathcal{D}$ –tangent surfaces, Darboux frame, Heisenberg group.

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### 1. Introduction

A smooth map  $\phi : N \rightarrow M$  is said to be biharmonic if it is a critical point of the bienergy functional:

$$E_2(\phi) = \int_N \frac{1}{2} |\mathcal{T}(\phi)|^2 dv_h,$$

where  $\mathcal{T}(\phi) := \text{tr} \nabla^\phi d\phi$  is the tension field of  $\phi$

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The Euler–Lagrange equation of the bienergy is given by  $\mathcal{T}_2(\phi) = 0$ . Here the section  $\mathcal{T}_2(\phi)$  is defined by

$$\mathcal{T}_2(\phi) = -\Delta_\phi \mathcal{T}(\phi) + \text{tr} R(\mathcal{T}(\phi), d\phi) d\phi,$$

and called the bitension field of  $\phi$ . Non-harmonic biharmonic maps are called proper biharmonic maps.

In this paper, we study  $\mathcal{D}$ –tangent surfaces of timelike biharmonic  $\mathcal{D}$ –helices according to Darboux frame on non-degenerate timelike surfaces in the Lorentzian Heisenberg group  $\mathbb{H}$ . We obtain parametric equation  $\mathcal{D}$ –tangent surfaces of timelike biharmonic  $\mathcal{D}$ –helices in the Lorentzian Heisenberg group  $\mathbb{H}$ . Moreover, we illustrate the figure of our main theorem.

## 2. Preliminaries

Heisenberg group plays an important role in many branches of mathematics such as representation theory, harmonic analysis, PDEs or even quantum mechanics, where it was initially defined as a group of  $3 \times 3$  matrices

$$\left\{ \begin{pmatrix} 1 & x & z \\ 0 & 1 & y \\ 0 & 0 & 1 \end{pmatrix} : x, y, z \in \mathbb{R} \right\}$$

with the usual multiplication rule.

The left-invariant Lorentz metric on  $\mathbb{H}$  is

$$\rho = -dx^2 + dy^2 + (xdy + dz)^2.$$

The following set of left-invariant vector fields forms an orthonormal basis for the corresponding Lie algebra:

$$\left\{ \mathbf{e}_1 = \frac{\partial}{\partial z}, \mathbf{e}_2 = \frac{\partial}{\partial y} - x \frac{\partial}{\partial z}, \mathbf{e}_3 = \frac{\partial}{\partial x} \right\}. \quad (2.1)$$

The characterising properties of this algebra are the following commutation relations:

$$\rho(\mathbf{e}_1, \mathbf{e}_1) = \rho(\mathbf{e}_2, \mathbf{e}_2) = 1, \quad \rho(\mathbf{e}_3, \mathbf{e}_3) = -1. \quad (2.2)$$

### 3. Timelike Biharmonic $\mathcal{D}$ -Helices According to Darboux Frame on a Non-Degenerate Timelike Surface in the Lorentzian Heisenberg Group $\mathbb{H}$

Let  $\Pi \subset \mathbb{H}$  be a timelike surface with the unit normal vector  $\mathbf{n}$  in the Lorentzian Heisenberg group  $\mathbb{H}$ . If  $\gamma$  is a timelike curve on  $\Pi \subset \mathbb{H}$ , then we have the Frenet frame  $\{\mathbf{T}, \mathbf{N}, \mathbf{B}\}$  and Darboux frame  $\{\mathbf{T}, \mathbf{n}, \mathbf{g}\}$  with spacelike vector  $\mathbf{g} = \mathbf{T} \wedge \mathbf{n}$  along the curve  $\gamma$ . Let  $\{\mathbf{T}, \mathbf{N}, \mathbf{B}\}$  be the Frenet frame fields tangent to the Lorentzian Heisenberg group  $\mathbb{H}$  along  $\gamma$  defined as follows:

$\mathbf{T}$  is the unit vector field  $\gamma'$  tangent to  $\gamma$ ,  $\mathbf{N}$  is the unit vector field in the direction of  $\nabla_{\mathbf{T}}\mathbf{T}$  (normal to  $\gamma$ ) and  $\mathbf{B}$  is chosen so that  $\{\mathbf{T}, \mathbf{N}, \mathbf{B}\}$  is a positively oriented orthonormal basis. Then, we have the following Frenet formulas:

$$\begin{aligned}\nabla_{\mathbf{T}}\mathbf{T} &= \kappa\mathbf{N}, \\ \nabla_{\mathbf{T}}\mathbf{N} &= \kappa\mathbf{T} + \tau\mathbf{B}, \\ \nabla_{\mathbf{T}}\mathbf{B} &= -\tau\mathbf{N},\end{aligned}\tag{3.1}$$

where  $\kappa$  is the curvature of  $\gamma$  and  $\tau$  is its torsion and

$$\begin{aligned}\rho(\mathbf{T}, \mathbf{T}) &= -1, \quad \rho(\mathbf{N}, \mathbf{N}) = 1, \quad \rho(\mathbf{B}, \mathbf{B}) = 1, \\ \rho(\mathbf{T}, \mathbf{N}) &= \rho(\mathbf{T}, \mathbf{B}) = \rho(\mathbf{N}, \mathbf{B}) = 0.\end{aligned}\tag{3.2}$$

Let  $\theta$  be the angle between  $\mathbf{N}$  and  $\mathbf{n}$ . The relationships between  $\{\mathbf{T}, \mathbf{N}, \mathbf{B}\}$  and  $\{\mathbf{T}, \mathbf{n}, \mathbf{g}\}$  are as follows:

$$\begin{aligned}\mathbf{T} &= \mathbf{T}, \\ \mathbf{N} &= \cos\theta\mathbf{n} + \sin\theta\mathbf{g}, \\ \mathbf{B} &= \sin\theta\mathbf{n} - \cos\theta\mathbf{g},\end{aligned}\tag{3.3}$$

and

$$\begin{aligned}\mathbf{T} &= \mathbf{T}, \\ \mathbf{g} &= \sin\theta\mathbf{N} - \cos\theta\mathbf{B}, \\ \mathbf{n} &= \cos\theta\mathbf{N} + \sin\theta\mathbf{B}.\end{aligned}\tag{3.4}$$

By differentiating (3.4), using (3.1), (3.3) and Frenet formulas we obtain

$$\begin{aligned}\nabla_{\mathbf{T}}\mathbf{T} &= (\kappa\cos\theta)\mathbf{n} + (\kappa\sin\theta)\mathbf{g}, \\ \nabla_{\mathbf{T}}\mathbf{g} &= (-\kappa\sin\theta)\mathbf{T} + \left(\tau + \frac{d\theta}{ds}\right)\mathbf{n}, \\ \nabla_{\mathbf{T}}\mathbf{n} &= (-\kappa\cos\theta)\mathbf{T} - \left(\tau + \frac{d\theta}{ds}\right)\mathbf{g}.\end{aligned}\tag{3.5}$$

If we represent  $\kappa \cos \theta$ ,  $\kappa \sin \theta$  and  $\tau + \frac{d\theta}{ds}$  with the symbols  $\kappa_n$ ,  $\kappa_g$ , and  $\tau_g$  respectively, then the equations in (3.5) can be written as

$$\begin{aligned}\nabla_{\mathbf{T}}\mathbf{T} &= \kappa_g\mathbf{g} + \kappa_n\mathbf{n}, \\ \nabla_{\mathbf{T}}\mathbf{g} &= -\kappa_g\mathbf{T} + \tau_g\mathbf{n}, \\ \nabla_{\mathbf{T}}\mathbf{n} &= -\kappa_n\mathbf{T} - \tau_g\mathbf{g}.\end{aligned}\tag{3.6}$$

With respect to the orthonormal basis  $\{\mathbf{e}_1, \mathbf{e}_2, \mathbf{e}_3\}$ , we can write

$$\begin{aligned}\mathbf{T} &= T_1\mathbf{e}_1 + T_2\mathbf{e}_2 + T_3\mathbf{e}_3, \\ \mathbf{g} &= g_1\mathbf{e}_1 + g_2\mathbf{e}_2 + g_3\mathbf{e}_3, \\ \mathbf{n} &= n_1\mathbf{e}_1 + n_2\mathbf{e}_2 + n_3\mathbf{e}_3,\end{aligned}\tag{3.7}$$

To separate a curve according to Darboux frame from that of Frenet- Serret frame, in the rest of the paper, we shall use notation for the curve as  $\mathcal{D}$ -curve.

First of all we recall the following well known result (cf. [8]).

**Theorem 3.1.** *Let  $\gamma : I \rightarrow \Pi \subset \mathbb{H}$  be a non-geodesic unit speed timelike curve on timelike surface  $\Pi$  in the Lorentzian Heisenberg group  $\mathbb{H}$ .  $\gamma$  is a unit speed timelike biharmonic curve on  $\Pi$  if and only if*

$$\begin{aligned}\kappa_n^2 + \kappa_g^2 &= \text{constant} \neq 0, \\ \kappa_n'' - \kappa_n^3 + \kappa_g\tau_g - \kappa_g^2\kappa_n + \kappa_g'\tau_g + \kappa_g\tau_g' - \tau_g^2\kappa_n &= \kappa_n(1 - 4g_1^2) - 4\kappa_g n_1 g_1, \\ \kappa_g'' - \kappa_g^3 - 2\kappa_n'\tau_g - \kappa_n^2\kappa_g - \kappa_n\tau_g' - \kappa_g\tau_g^2 &= 4\kappa_n n_1 g_1 + \kappa_g(1 - 4n_1^2).\end{aligned}\tag{3.8}$$

**Theorem 3.2.** *Let  $\gamma : I \rightarrow \Pi \subset \mathbb{H}$  be a non-geodesic unit speed timelike biharmonic  $\mathcal{D}$ - helix on timelike surface  $\Pi$  in the Lorentzian Heisenberg group  $\mathbb{H}$ . Then parametric equations of timelike biharmonic  $\mathcal{D}$ - helix are*

$$\begin{aligned}\mathbf{x}(s) &= \frac{\cosh \aleph}{\mathcal{R}}(\cosh[\mathcal{R}s] \sinh[\wp] + \cosh[\wp] \sinh[\mathcal{R}s]) + \wp_1, \\ \mathbf{y}(s) &= \frac{\cosh \aleph}{\mathcal{R}}(\cosh[\wp] \cosh[\mathcal{R}s] + \sinh[\wp] \sinh[\mathcal{R}s]) + \wp_2, \\ \mathbf{z}(s) &= \sinh \aleph s + \frac{(\wp + \mathcal{R}s)}{2\mathcal{R}^2} \cosh^2 \aleph - \frac{\wp_1}{\mathcal{R}} \cosh^2 \aleph \cosh[\wp] \cosh[\mathcal{R}s] \\ &\quad - \frac{\wp_1}{\mathcal{R}} \cosh^2 \aleph \sinh[\wp] \sinh[\mathcal{R}s] - \frac{1}{4\mathcal{R}^2} \cosh^2 \aleph \sinh 2[\mathcal{R}s + \wp] + \wp_3,\end{aligned}\tag{3.9}$$

where  $\wp, \wp_1, \wp_2, \wp_3$ , are constants of integration and

$$\mathcal{R} = \text{sech } \aleph \sqrt{\kappa_n^2 + \kappa_g^2} - 2 \sinh \aleph.$$

Using Mathematica in above system we have Fig.1:

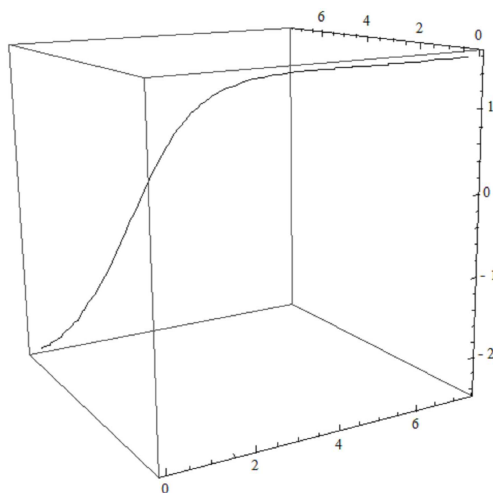


Fig. 1

**4.  $\mathcal{D}$ -Tangent Surfaces According to Darboux Frame on a Non-Degenerate Timelike Surface in the Lorentzian Heisenberg Group  $\mathbb{H}$**

A ruled surface is formed by a continuous family of straight line segments.

The  $\mathcal{D}$ -tangent surface of  $\gamma$  is a ruled surface

$$\mathcal{O}(s, u) = \gamma(s) + u\mathbf{T}(s). \tag{4.1}$$

**Theorem 4.1.** *Let  $\gamma : I \rightarrow \mathcal{O} \subset \mathbb{H}$  be a non-geodesic unit speed timelike biharmonic  $\mathcal{D}$ -helix on timelike surface  $\mathcal{O}$  in the Lorentzian Heisenberg group  $\mathbb{H}$ . Then,*

equation of  $\mathcal{D}$ -tangent surface of timelike biharmonic  $\mathcal{D}$ - helix is

$$\begin{aligned}
\mathcal{O}(s, u) = & \left[ \sinh \aleph s + \frac{(\wp + \mathcal{R}s)}{2\mathcal{R}^2} \cosh^2 \aleph - \frac{\wp_1}{\mathcal{R}} \cosh^2 \aleph \cosh[\wp] \cosh[\mathcal{R}s] \right. \\
& - \frac{\wp_1}{\mathcal{R}} \cosh^2 \aleph \sinh[\wp] \sinh[\mathcal{R}s] - \frac{1}{4\mathcal{R}^2} \cosh^2 \aleph \sinh 2[\mathcal{R}s + \wp] \\
& + \wp_3 + \left[ \frac{\cosh \aleph}{\mathcal{R}} (\cosh[\mathcal{R}s] \sinh[\wp] + \cosh[\wp] \sinh[\mathcal{R}s]) + \wp_1 \right] \quad (4.2) \\
& \left[ \frac{\cosh \aleph}{\mathcal{R}} (\cosh[\wp] \cosh[\mathcal{R}s] + \sinh[\wp] \sinh[\mathcal{R}s]) + \wp_2 \right] + u \sinh \mathcal{P} \mathbf{e}_1 \\
& + \left[ \frac{\cosh \aleph}{\mathcal{R}} (\cosh[\wp] \cosh[\mathcal{R}s] + \sinh[\wp] \sinh[\mathcal{R}s]) \right. \\
& + u \cosh \mathcal{P} \sinh [\mathcal{M}s + \mathcal{N}] + \wp_2 \left. \right] \mathbf{e}_2 \\
& + \left[ \frac{\cosh \aleph}{\mathcal{R}} (\cosh[\mathcal{R}s] \sinh[\wp] + \cosh[\wp] \sinh[\mathcal{R}s]) \right. \\
& \left. + u \cosh \mathcal{P} \cosh [\mathcal{M}s + \mathcal{N}] + \wp_1 \right] \mathbf{e}_3,
\end{aligned}$$

where  $\wp, \wp_1, \wp_2, \wp_3$ , are constants of integration and

$$\mathcal{R} = \operatorname{sech} \aleph \sqrt{\kappa_n^2 + \kappa_g^2} - 2 \sinh \aleph.$$

**Proof:** From the assumption we get

$$\mathbf{T} = \sinh \mathcal{P} \mathbf{e}_1 + \cosh \mathcal{P} \sinh [\mathcal{M}s + \mathcal{N}] \mathbf{e}_2 + \cosh \mathcal{P} \cosh [\mathcal{M}s + \mathcal{N}] \mathbf{e}_3. \quad (4.3)$$

Substituting Eq.(3.9) and Eq.(4.3) into Eq.(4.1) , we obtain Eq.(4.2). This completes the proof.  $\square$

**Theorem 4.2.** *Let  $\gamma : I \rightarrow \mathcal{O} \subset \mathbb{H}$  be a non-geodesic unit speed timelike biharmonic  $\mathcal{D}$ - helix on timelike surface  $\mathcal{O}$  in the Lorentzian Heisenberg group  $\mathbb{H}$ . Then, equation of  $\mathcal{D}$ -tangent surface of timelike biharmonic  $\mathcal{D}$ - helix are*

$$\begin{aligned}
x_{\mathcal{O}} = & \left[ \frac{\cosh \aleph}{\mathcal{R}} (\cosh[\mathcal{R}s] \sinh[\wp] + \cosh[\wp] \sinh[\mathcal{R}s]) + u \cosh \mathcal{P} \cosh [\mathcal{M}s + \mathcal{N}] + \wp_1 \right] \\
y_{\mathcal{O}} = & \left[ \frac{\cosh \aleph}{\mathcal{R}} (\cosh[\wp] \cosh[\mathcal{R}s] + \sinh[\wp] \sinh[\mathcal{R}s]) + u \cosh \mathcal{P} \sinh [\mathcal{M}s + \mathcal{N}] + \wp_2 \right] \\
z_{\mathcal{O}} = & \left[ \sinh \aleph s + \frac{(\wp + \mathcal{R}s)}{2\mathcal{R}^2} \cosh^2 \aleph - \frac{\wp_1}{\mathcal{R}} \cosh^2 \aleph \cosh[\wp] \cosh[\mathcal{R}s] \right. \\
& - \frac{\wp_1}{\mathcal{R}} \cosh^2 \aleph \sinh[\wp] \sinh[\mathcal{R}s] - \frac{1}{4\mathcal{R}^2} \cosh^2 \aleph \sinh 2[\mathcal{R}s + \wp] \\
& + \wp_3 + \left[ \frac{\cosh \aleph}{\mathcal{R}} (\cosh[\mathcal{R}s] \sinh[\wp] + \cosh[\wp] \sinh[\mathcal{R}s]) + \wp_1 \right] \\
& \left[ \frac{\cosh \aleph}{\mathcal{R}} (\cosh[\wp] \cosh[\mathcal{R}s] + \sinh[\wp] \sinh[\mathcal{R}s]) + \wp_2 \right] + u \sinh \mathcal{P} \\
& - \left[ \frac{\cosh \aleph}{\mathcal{R}} (\cosh[\wp] \cosh[\mathcal{R}s] + \sinh[\wp] \sinh[\mathcal{R}s]) + u \cosh \mathcal{P} \sinh [\mathcal{M}s + \mathcal{N}] + \wp_2 \right] \\
& \left. \left[ \frac{\cosh \aleph}{\mathcal{R}} (\cosh[\mathcal{R}s] \sinh[\wp] + \cosh[\wp] \sinh[\mathcal{R}s]) + u \cosh \mathcal{P} \cosh [\mathcal{M}s + \mathcal{N}] + \wp_1 \right], \right.
\end{aligned}$$

where  $\wp, \wp_1, \wp_2, \wp_3$ , are constants of integration and

$$\mathcal{R} = \operatorname{sech} \aleph \sqrt{\kappa_n^2 + \kappa_g^2} - 2 \sinh \aleph.$$

**Proof:** From Theorem 4.1, we easily have above system, which completes the proof.  $\square$

In the light of Theorem 4.2, we give the following figures for the  $\mathcal{D}$ -tangent surface of timelike biharmonic  $\mathcal{D}$ -helix.

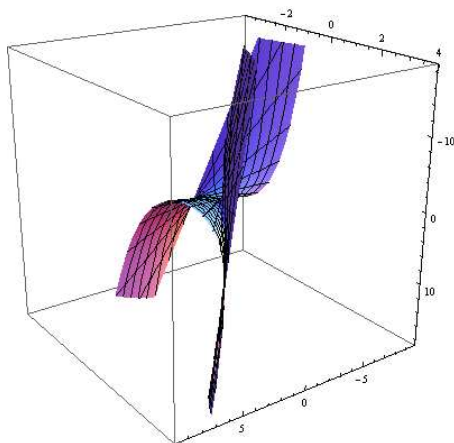


Fig. 2

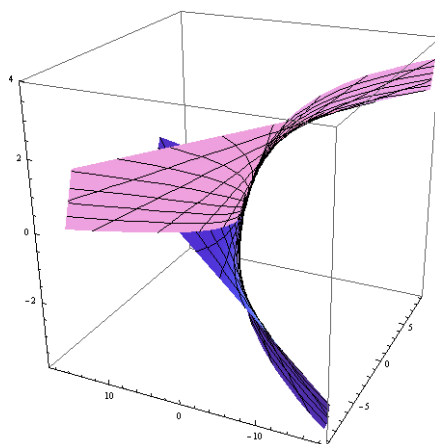


Fig. 3

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