



Biharmonic \mathcal{S} -Curves According to Sabban Frame in Heisenberg Group Heis^3

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ABSTRACT: In this paper, we study biharmonic curves according to Sabban frame in the Heisenberg group Heis^3 . We characterize the biharmonic curves in terms of their geodesic curvature and we prove that all of biharmonic curves are helices in the Heisenberg group Heis^3 . Finally, we find out their explicit parametric equations according to Sabban Frame.

Key Words: Biharmonic curve, Heisenberg group, curvature, torsion.

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1. Introduction

Harmonic maps $f : (M, g) \rightarrow (N, h)$ between manifolds are the critical points of the energy

$$E(f) = \frac{1}{2} \int_M e(f) v_g, \quad (1.1)$$

where v_g is the volume form on (M, g) and

$$e(f)(x) := \frac{1}{2} \|df(x)\|_{T^*M \otimes f^{-1}TN}^2$$

is the energy density of f at the point $x \in M$.

Critical points of the energy functional are called harmonic maps.

The first variational formula of the energy gives the following characterization of harmonic maps: the map f is harmonic if and only if its tension field $\tau(f)$ vanishes identically, where the tension field is given by

$$\tau(f) = \text{trace} \nabla df. \quad (1.2)$$

As suggested by Eells and Sampson in [6], we can define the bienergy of a map f by

$$E_2(f) = \frac{1}{2} \int_M \|\tau(f)\|^2 v_g, \quad (1.3)$$

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and say that is biharmonic if it is a critical point of the bienergy.

Jiang derived the first and the second variation formula for the bienergy in [7,8], showing that the Euler–Lagrange equation associated to E_2 is

$$\begin{aligned}\tau_2(f) &= -\mathcal{J}^f(\tau(f)) = -\Delta\tau(f) - \text{trace}R^N(df, \tau(f))df \\ &= 0,\end{aligned}\quad (1.4)$$

where \mathcal{J}^f is the Jacobi operator of f . The equation $\tau_2(f) = 0$ is called the biharmonic equation. Since \mathcal{J}^f is linear, any harmonic map is biharmonic. Therefore, we are interested in proper biharmonic maps, that is non-harmonic biharmonic maps.

This study is organised as follows: Firstly, we study biharmonic curves according to Sabban frame in the Heisenberg group Heis^3 . Secondly, we characterize the biharmonic curves in terms of their geodesic curvature and we prove that all of biharmonic curves are helices in the Heisenberg group Heis^3 . Finally, we find out their explicit parametric equations according to Sabban Frame.

2. The Heisenberg Group Heis^3

Heisenberg group Heis^3 can be seen as the space \mathbb{R}^3 endowed with the following multiplication:

$$(\bar{x}, \bar{y}, \bar{z})(x, y, z) = (\bar{x} + x, \bar{y} + y, \bar{z} + z - \frac{1}{2}\bar{x}y + \frac{1}{2}x\bar{y}) \quad (2.1)$$

Heis^3 is a three-dimensional, connected, simply connected and 2-step nilpotent Lie group.

The Riemannian metric g is given by

$$g = dx^2 + dy^2 + (dz - xdy)^2.$$

The Lie algebra of Heis^3 has an orthonormal basis

$$\mathbf{e}_1 = \frac{\partial}{\partial x}, \quad \mathbf{e}_2 = \frac{\partial}{\partial y} + x\frac{\partial}{\partial z}, \quad \mathbf{e}_3 = \frac{\partial}{\partial z}, \quad (2.2)$$

for which we have the Lie products

$$[\mathbf{e}_1, \mathbf{e}_2] = \mathbf{e}_3, \quad [\mathbf{e}_2, \mathbf{e}_3] = [\mathbf{e}_3, \mathbf{e}_1] = 0$$

with

$$g(\mathbf{e}_1, \mathbf{e}_1) = g(\mathbf{e}_2, \mathbf{e}_2) = g(\mathbf{e}_3, \mathbf{e}_3) = 1.$$

We obtain

$$\begin{aligned}\nabla_{\mathbf{e}_1}\mathbf{e}_1 &= \nabla_{\mathbf{e}_2}\mathbf{e}_2 = \nabla_{\mathbf{e}_3}\mathbf{e}_3 = 0, \\ \nabla_{\mathbf{e}_1}\mathbf{e}_2 &= -\nabla_{\mathbf{e}_2}\mathbf{e}_1 = \frac{1}{2}\mathbf{e}_3, \\ \nabla_{\mathbf{e}_1}\mathbf{e}_3 &= \nabla_{\mathbf{e}_3}\mathbf{e}_1 = -\frac{1}{2}\mathbf{e}_2, \\ \nabla_{\mathbf{e}_2}\mathbf{e}_3 &= \nabla_{\mathbf{e}_3}\mathbf{e}_2 = \frac{1}{2}\mathbf{e}_1.\end{aligned}$$

The components $\{R_{ijkl}\}$ of R relative to $\{\mathbf{e}_1, \mathbf{e}_2, \mathbf{e}_3\}$ are defined by

$$R_{ijk} = R(\mathbf{e}_i, \mathbf{e}_j)\mathbf{e}_k, \quad R_{ijkl} = R(\mathbf{e}_i, \mathbf{e}_j, \mathbf{e}_k, \mathbf{e}_l) = g(R(\mathbf{e}_i, \mathbf{e}_j)\mathbf{e}_l, \mathbf{e}_k).$$

The non vanishing components of the above tensor fields are

$$R_{121} = \frac{3}{4}\mathbf{e}_2, \quad R_{131} = -\frac{1}{4}\mathbf{e}_3, \quad R_{122} = -\frac{3}{4}\mathbf{e}_1,$$

$$R_{232} = -\frac{1}{4}\mathbf{e}_3, \quad R_{133} = \frac{1}{4}\mathbf{e}_1, \quad R_{233} = \frac{1}{4}\mathbf{e}_2,$$

and

$$R_{1212} = -\frac{3}{4}, \quad R_{1313} = R_{2323} = \frac{1}{4}. \quad (2.3)$$

3. Biharmonic S-Curves According To Sabban Frame In The Heisenberg Group $Heis^3$

Let $\gamma : I \rightarrow Heis^3$ be a non geodesic curve on the Heisenberg group $Heis^3$ parametrized by arc length. Let $\{\mathbf{T}, \mathbf{N}, \mathbf{B}\}$ be the Frenet frame fields tangent to the Heisenberg group $Heis^3$ along γ defined as follows:

\mathbf{T} is the unit vector field γ' tangent to γ , \mathbf{N} is the unit vector field in the direction of $\nabla_{\mathbf{T}}\mathbf{T}$ (normal to γ), and \mathbf{B} is chosen so that $\{\mathbf{T}, \mathbf{N}, \mathbf{B}\}$ is a positively oriented orthonormal basis. Then, we have the following Frenet formulas:

$$\begin{aligned} \nabla_{\mathbf{T}}\mathbf{T} &= \kappa\mathbf{N}, \\ \nabla_{\mathbf{T}}\mathbf{N} &= -\kappa\mathbf{T} + \tau\mathbf{B}, \\ \nabla_{\mathbf{T}}\mathbf{B} &= -\tau\mathbf{N}, \end{aligned} \quad (3.1)$$

where κ is the curvature of γ and τ is its torsion,

$$\begin{aligned} g(\mathbf{T}, \mathbf{T}) &= 1, \quad g(\mathbf{N}, \mathbf{N}) = 1, \quad g(\mathbf{B}, \mathbf{B}) = 1, \\ g(\mathbf{T}, \mathbf{N}) &= g(\mathbf{T}, \mathbf{B}) = g(\mathbf{N}, \mathbf{B}) = 0. \end{aligned}$$

Now we give a new frame different from Frenet frame. Let $\alpha : I \rightarrow \mathbb{S}_{Heis^3}^2$ be unit speed spherical curve. We denote σ as the arc-length parameter of α . Let us denote $\mathbf{t}(\sigma) = \alpha'(\sigma)$, and we call $\mathbf{t}(\sigma)$ a unit tangent vector of α . We now set a vector $\mathbf{s}(\sigma) = \alpha(\sigma) \times \mathbf{t}(\sigma)$ along α . This frame is called the Sabban frame of α on the Heisenberg group $Heis^3$. Then we have the following spherical Frenet-Serret formulae of α :

$$\begin{aligned} \nabla_{\mathbf{t}}\alpha &= \mathbf{t}, \\ \nabla_{\mathbf{t}}\mathbf{t} &= -\alpha + \kappa_g\mathbf{s}, \\ \nabla_{\mathbf{t}}\mathbf{s} &= -\kappa_g\mathbf{t}, \end{aligned} \quad (3.2)$$

where κ_g is the geodesic curvature of the curve α on the $\mathbb{S}_{Heis^3}^2$ and

$$\begin{aligned} g(\mathbf{t}, \mathbf{t}) &= 1, \quad g(\alpha, \alpha) = 1, \quad g(\mathbf{s}, \mathbf{s}) = 1, \\ g(\mathbf{t}, \alpha) &= g(\mathbf{t}, \mathbf{s}) = g(\alpha, \mathbf{s}) = 0. \end{aligned}$$

With respect to the orthonormal basis $\{\mathbf{e}_1, \mathbf{e}_2, \mathbf{e}_3\}$, we can write

$$\begin{aligned} \alpha &= \alpha_1 \mathbf{e}_1 + \alpha_2 \mathbf{e}_2 + \alpha_3 \mathbf{e}_3, \\ \mathbf{t} &= t_1 \mathbf{e}_1 + t_2 \mathbf{e}_2 + t_3 \mathbf{e}_3, \\ \mathbf{s} &= s_1 \mathbf{e}_1 + s_2 \mathbf{e}_2 + s_3 \mathbf{e}_3. \end{aligned} \tag{3.3}$$

To separate a biharmonic curve according to Sabban frame from that of Frenet-Serret frame, in the rest of the paper, we shall use notation for the curve defined above as biharmonic \mathcal{S} -curve.

Theorem 3.1. $\alpha : I \longrightarrow \mathbb{S}_{Heis^3}^2$ is a biharmonic \mathcal{S} -curve if and only if

$$\begin{aligned} \kappa_g &= \text{constant} \neq 0, \\ 1 + \kappa_g^2 &= -\left[\frac{1}{4} - s_3^2\right] + \kappa_g[-\alpha_3 s_3], \\ \kappa_g'' - \kappa_g^3 &= \alpha_3 s_3 + \kappa_g\left[\frac{1}{4} - \alpha_3^2\right]. \end{aligned} \tag{3.4}$$

Proof: Using (2.1) and Sabban formulas (3.2), we have (3.4). \square

Corollary 3.2. $\alpha : I \longrightarrow \mathbb{S}_{Heis^3}^2$ is a biharmonic \mathcal{S} -curve if and only if

$$\begin{aligned} \kappa_g &= \text{constant} \neq 0, \\ 1 + \kappa_g^2 &= -\left[\frac{1}{4} - s_3^2\right] + \kappa_g[-\alpha_3 s_3], \\ \kappa_g^3 &= -\alpha_3 s_3 - \kappa_g\left[\frac{1}{4} - \alpha_3^2\right]. \end{aligned} \tag{3.4}$$

Lemma 3.3. All of biharmonic \mathcal{S} -curves in $\mathbb{S}_{Heis^3}^2$ are helices.

Theorem 3.4. Let $\alpha : I \longrightarrow \mathbb{S}_{Heis^3}^2$ be a unit speed non-geodesic biharmonic

S-curve. Then, the parametric equations of α are

$$\begin{aligned} x^S(\sigma) &= -\frac{\sin^2 \varphi}{(\sqrt{1+\kappa_g^2} - \sin \varphi \cos \varphi)} \cos\left[\left(\frac{\sqrt{1+\kappa_g^2}}{\sin \varphi} - \cos \varphi\right)\sigma + \mathcal{M}_1\right] + \mathcal{M}_2, \\ y^S(\sigma) &= \frac{\sin^2 \varphi}{(\sqrt{1+\kappa_g^2} - \sin \varphi \cos \varphi)} \sin\left[\left(\frac{\sqrt{1+\kappa_g^2}}{\sin \varphi} - \cos \varphi\right)\sigma + \mathcal{M}_1\right] + \mathcal{M}_3, \quad (3.5) \\ z^S(\sigma) &= \cos \varphi \sigma - \sin \varphi \frac{(\sqrt{1+\kappa_g^2} - \sin \varphi \cos \varphi)\sigma + \mathcal{M}_1}{2(\sqrt{1+\kappa_g^2} - \sin \varphi \cos \varphi)} \\ &\quad - \sin^2 \varphi \frac{\sin 2\left[\left(\frac{\sqrt{1+\kappa_g^2}}{\sin \varphi} - \cos \varphi\right)\sigma + \mathcal{M}_1\right]}{4(\sqrt{1+\kappa_g^2} - \sin \varphi \cos \varphi)} \\ &\quad + \frac{\mathcal{M}_2}{(\sqrt{1+\kappa_g^2} - \sin \varphi \cos \varphi)} \sin^3 \varphi \sin\left[\left(\frac{\sqrt{1+\kappa_g^2}}{\sin \varphi} - \cos \varphi\right)\sigma + \mathcal{M}_1\right] + \mathcal{M}_4, \end{aligned}$$

where $\mathcal{M}_1, \mathcal{M}_2, \mathcal{M}_3, \mathcal{M}_4$ are constants of integration.

Proof: Since α is biharmonic, α is a S-helix. So, without loss of generality, we take the axis of α is parallel to the vector \mathbf{e}_3 . Then,

$$g(\mathbf{t}, \mathbf{e}_3) = t_3 = \cos \varphi, \quad (3.6)$$

where φ is constant angle.

So, substituting the components t_1, t_2 and t_3 in the equation (3.3), we have the following equation

$$\mathbf{t} = \sin \varphi \sin \mu \mathbf{e}_1 + \sin \varphi \cos \mu \mathbf{e}_2 + \cos \varphi \mathbf{e}_3. \quad (3.7)$$

The covariant derivative of the vector field t is:

$$\nabla_{\mathbf{t}} \mathbf{t} = (t'_1 + t_2 t_3) \mathbf{e}_1 + (t'_2 - t_1 t_3) \mathbf{e}_2 + t'_3 \mathbf{e}_3. \quad (3.8)$$

From above equation we have

$$\mu(\sigma) = \left(\frac{\sqrt{1+\kappa_g^2}}{\sin \varphi} - \cos \varphi\right)\sigma + \mathcal{M}_1, \quad (3.10)$$

where \mathcal{M}_1 is a constant of integration.

Thus (3.9) and (3.10), imply

$$\begin{aligned} \mathbf{t} &= \sin \varphi \sin\left[\left(\frac{\sqrt{1+\kappa_g^2}}{\sin \varphi} - \cos \varphi\right)\sigma + \mathcal{M}_1\right] \mathbf{e}_1 \\ &\quad + \sin \varphi \cos\left[\left(\frac{\sqrt{1+\kappa_g^2}}{\sin \varphi} - \cos \varphi\right)\sigma + \mathcal{M}_1\right] \mathbf{e}_2 + \cos \varphi \mathbf{e}_3. \end{aligned} \quad (3.11)$$

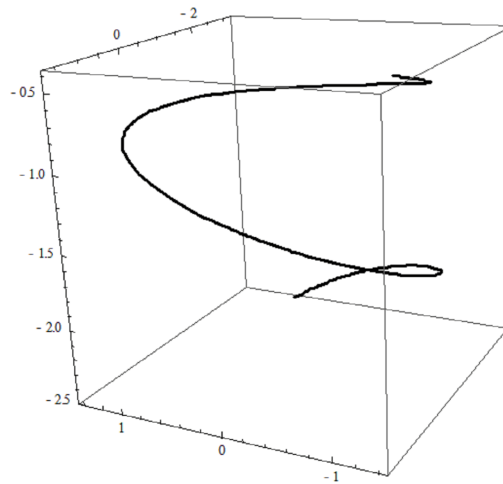
Using (2.1) in (3.11), we obtain

$$\begin{aligned} \mathbf{t} = & \left(\sin \varphi \sin \left[\left(\frac{\sqrt{1 + \kappa_g^2}}{\sin \varphi} - \cos \varphi \right) \sigma + \mathcal{M}_1 \right], \sin \varphi \cos \left[\left(\frac{\sqrt{1 + \kappa_g^2}}{\sin \varphi} - \cos \varphi \right) \sigma + \mathcal{M}_1 \right], \right. \\ & \left. \cos \varphi + \sin \varphi \left(- \frac{\sin^2 \varphi}{\left(\sqrt{1 + \kappa_g^2} - \sin \varphi \cos \varphi \right)} \cos \left[\left(\frac{\sqrt{1 + \kappa_g^2}}{\sin \varphi} - \cos \varphi \right) \sigma + \mathcal{M}_1 \right] \right) \right. \\ & \left. + \mathcal{M}_2 \right) \cos \left[\left(\frac{\sqrt{1 + \kappa_g^2}}{\sin \varphi} - \cos \varphi \right) \sigma + \mathcal{M}_1 \right]. \end{aligned} \quad (3.12)$$

where $\mathcal{M}_1, \mathcal{M}_2$ are constants of integration.

Integrating both sides, we have (3.9). This proves our assertion. Thus, the proof of theorem is completed. \square

We can use Mathematica in above theorem, yields



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