



Stabilization of Navier-Stokes Equations

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ABSTRACT: We survey here a few recent stabilization results for Navier-Stokes equations.

Key Words: Navier-Stokes equations, Hamilton-Jacobi equation, Brownian motion.

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1. Introduction

We are concerned here with the stabilization of steady state solutions to Navier-Stokes equation

$$\begin{aligned} y_t - \nu \Delta y + (y \cdot \nabla)y &= \nabla p + mu + f_e && \text{in } D \times R^+ \\ \nabla \cdot y &= 0 && \text{in } D \times R^+ \\ y &= 0 && \text{on } \partial D \times R^+ \\ y(0) &= y_0 && \text{in } D \end{aligned} \quad (1)$$

where D is a bounded and smooth domain of R^d , $d = 2, 3$ and $f_e \in (L^2(D))^d$. Here $m \in C_0^1(\mathcal{O})$ and $m > 1$ on $\mathcal{O}_1 \subset \subset \mathcal{O}$ where $\mathcal{O}_1, \mathcal{O}$ are open subdomains of D .

There is a large number of recent works devoted to feedback stabilization of Navier-Stokes equations of the form (1) with internal and boundary controllers.

Let y_e be a solution to stationary equation

$$\begin{aligned} -\nu \Delta y_e + (y_e \cdot \nabla)y_e &= \nabla p_e + f_e && \text{in } D \\ \nabla \cdot y_e &= 0 && \text{in } D \\ y_e &= 0 && \text{on } \partial D. \end{aligned} \quad (2)$$

The main result established in [4] (see also [1], [5], [6]) is that *there is a stabilizing feedback controller $u = u(t, x)$ of the form*

$$u(t, x) = - \sum_{i=1}^M \psi_i(x) \int_D R(y - y_e) \psi_i m dx \quad (3)$$

where R is a self-adjoint operator to be precised below and $\{\psi_i\}_{i=1}^M$ is a given system of functions.

We set $y - y_e \implies y$, $H = \{y \in (L^2(D))^d; \nabla \cdot y = 0, y \cdot n = 0 \text{ on } \partial D\}$, $P : (L^2(D))^d \rightarrow H$ is the Leray projector and

$$Ay = P\Delta y, \quad By = P(y \cdot \nabla)y,$$

with $D(A) = \{y \in H \cap (H_0^1(D))^d; \Delta y \in (L^2(D))^d\}$. Then we may write equation for $y \implies y - y_e$ as

$$\begin{aligned} \frac{dy}{dt}(t) + \nu Ay(t) + A_0 y(t) + By(t) &= Fu \\ y(0) &= y_0 - y_e = x \end{aligned} \quad (4)$$

where $A_0 y = P((y_e \cdot \nabla)y + (y \cdot \nabla)y_e)$ and $Fu = \sum_{i=1}^M P(m\psi_i)u_i$, $\{\psi_i\}_{i=1}^M \subset D(A)$. We shall denote by $|\cdot|$ the norm of H , (\cdot, \cdot) the scalar product and by A^α , $0 < \alpha < 1$, the fractional power of order α of operator A^α . We set also $|u|_\alpha = |A^{\frac{\alpha}{2}}u|$ for all $\alpha \in (0, 1)$.

The operator R arising in (3) is the symmetric solution to Riccati equation

$$((\nu A + A_0)y, Ry) + \frac{1}{2}|F^* Ry|^2 = \frac{1}{2}|A^{\frac{3}{4}}y|^2, \quad \forall y \in D(A^{\frac{3}{4}}) \quad (5)$$

and has the following properties (see [4])

$$R \in L(D(A^{\frac{1}{4}}), (D(A^{\frac{1}{4}}))' \cap L(D(A^{\frac{1}{2}}), H)) \quad (6)$$

$$(Ry, y) \geq \delta |A^{\frac{1}{4}}y|^2, \quad \forall y \in D(A^{\frac{1}{4}}). \quad (7)$$

Here $\{\psi_i\}_{i=1}^M$ is a system of functions which belongs to space $\text{lin}\{\varphi_j\}_{j=1}^N$ of eigenfunctions for $\mathcal{A} = \nu A + A_0$ (N is the number of unstable eigenvalues) and the dimension M of the system is given by spectral properties of \mathcal{A} ($M = 2$ if all unstable eigenvalues are simple and is maximum $2N$ in general case). As shown in [4], the feedback operator (3) exponentially stabilizes zero solution of (4) in a neighborhood

$$\mathcal{U}_\rho = \{x \in D(A^{\frac{1}{4}}); |x|_{\frac{1}{2}} < \rho\}. \quad (8)$$

The optimal radius ρ of stability domain \mathcal{U}_ρ is determined by formula

$$\max_{|y|_{\frac{1}{2}} \leq \rho} \frac{2|b(y, y, Ry)|}{|y|_{\frac{3}{2}}^2} < 1 \quad (9)$$

where $b(y, z, w) = \int_D (y \cdot \nabla)zw \, dx$, $\forall y, z, w \in D(A^{\frac{1}{2}})$. It follows that for $d = 3$, ρ might be taken as

$$0 < \rho < \frac{1}{2} \|R\|_{L(D(A^{\frac{1}{2}}), H)}. \quad (10)$$

This result is of course not optimal because it uses a linear stabilizing feedback for the linearized equation in nonlinear equation (4) and so our aim here is to find the maximal domain of stability via a nonlinear feedback control.

One might speculate that an optimal control feedback for equation (4) with cost functional

$$J(y, u) = \frac{1}{2} \int_0^\infty (|y(t)|_{\frac{3}{2}}^2 + |u(t)|_M^2) dt \quad (11)$$

has a wider domain of stability. The construction of such a feedback law is our aim. Here $|\cdot|_M$ is the norm in R^M .

We shall denote by $\mathcal{A} : D(\mathcal{A}) = D(A) \rightarrow H$ the operator

$$\mathcal{A}y = \nu Ay + A_0y, \quad D(\mathcal{A}) = D(A) \quad (12)$$

and by $B : D(B) \subset H \rightarrow H$ the operator defined by

$$(By, w) = b(y, y, w), \quad \forall w \in D(A^{\frac{1}{2}}). \quad (13)$$

We shall denote by W the space $D(A^{\frac{1}{4}})$ with the norm denoted $|\cdot|_{\frac{1}{2}}$.

In Section 2 below we shall develop this approach and refer to [2] for complete proofs. Section 3 is devoted to a different technique, stabilization by noise developed in forth-coming paper [3].

2. The nonlinear optimal control problem

Consider the control system (4), i.e.,

$$\begin{aligned} \frac{dy}{dt} + \mathcal{A}y + By &= Fu, \quad \forall t \geq 0, \\ y(0) &= x, \end{aligned} \quad (14)$$

where $u(t) = \{u_i(t)\}_{i=1}^M$ and $Fu = \sum_{i=1}^M u_i P(m\psi_i)(x)$. By solution to (14) on $[0, T]$ we

mean a function $y \in C([0, T]; H) \cap L^2(0, T; D(A^{\frac{1}{2}}))$ with $\frac{dy}{dt} \in L^2(0, T; (D(A^{\frac{1}{2}}))')$ which satisfies a.e. the equation. For $d = 2$ and $x \in H$ there is a unique such a solution while for $d = 3$ it exists only locally or globally (in time) for x in a suitable chosen neighborhood of origin. We shall denote by \mathcal{D} the stabilizability domain of (14) with respect to cost functional (11), i.e.,

$$\mathcal{D} = \{x \in W; \exists (y, u) \in L^2(0, \infty; D(A^{\frac{3}{4}})) \times L^2(0, \infty; R^M) \text{ satisfying (14)}\}. \quad (15)$$

As mentioned earlier $\mathcal{D} \neq \emptyset$ and it contains \mathcal{U}_ρ given by (8).

Define the function $\varphi : \mathcal{D} \rightarrow R$,

$$\varphi(x) = \inf_{(y, u)} \{J(y, u)\}, \quad \forall x \in \mathcal{D}. \quad (16)$$

We have

Proposition 1 For each $x \in \mathcal{D}$ there is at least one pair (y^*, u^*) such that

$$\varphi(x) = J(y^*, u^*). \quad (17)$$

Moreover, $y^*(t) \in \mathcal{D}$, $\forall t \geq 0$.

Proof. Existence is standard and so it will be omitted. \blacksquare

Theorem 2 below is a maximum principle type result for problem (16). For the sake of simplicity we shall assume from now on that $d = 2$. The extension to $d = 3$ is however straightforward.

Theorem 2 Let (y^*, u^*) be optimal in problem (16). Then

$$u^*(t) = F^*p(t) = \left\{ \int_D m(x)\psi_i(x)p(t, x)dx \right\}_{i=1}^M, \quad \forall t > 0, \quad (18)$$

where $p \in L^2(0, \infty; H) \cap C([0, \infty); H) \cap L^\infty(0, \infty; D(A^{\frac{1}{4}})) \cap L^2(0, \infty; D(A^{\frac{3}{4}}))$ is the solution to equation

$$\frac{dp}{dt} - \mathcal{A}^*p - (B'(y^*))^*p = A^{\frac{3}{2}}y^* \quad a.e. \quad t \geq 0. \quad (19)$$

Here \mathcal{A}^* is the adjoint of \mathcal{A} in H and $(B'(y^*))^*$ is defined by

$$((B'(y^*))^*p, w) = b(y^*, w, p) + b(w, y^*, p), \quad \forall w \in D(A^{\frac{1}{2}}). \quad (20)$$

Proof: If (y^*, u^*) is optimal in (16), then it is also optimal for problem

$$\begin{aligned} & \text{Min} \left\{ \int_0^\infty (|y(t)|_{\frac{2}{3}}^2 + |v(t) - F^*Ry(t)|^2) dt; \right. \\ & \left. \frac{dy}{dt} + \mathcal{A} + FF^*Ry + By = Fv, \quad y(0) = x, \quad v \in L^2(0, \infty; R^M) \right\}, \end{aligned} \quad (21)$$

too, where $R \in L(D(A^{\frac{1}{4}}), (D(A^{\frac{1}{4}})) \cap L(D(A^{\frac{1}{2}}), H))$ is the solution to the algebraic Riccati equation (5).

Next we consider the operator $\mathcal{L} : L^2(0, \infty; H) \rightarrow L^2(0, \infty; H)$ defined by

$$(\mathcal{L}z)(t) = \frac{dz}{dt} + \mathcal{A}z(t) + B'((y^*(t))z(t) + FF^*Rz(t)), \quad \forall z \in D(\mathcal{L}) \quad (22)$$

$$\begin{aligned} D(\mathcal{L}) = & \left\{ z \in L^2(0, \infty; D(A^{\frac{3}{4}})) \cap C([0, \infty); D(A^{\frac{1}{4}})); \right. \\ & \frac{dz}{dt} \in L^2(0, \infty; (D(A^{\frac{3}{4}}))'), \\ & \left. \frac{dz}{dt} + \mathcal{A}z \in L^2(0, \infty; H), \quad z(0) = 0 \right\}. \end{aligned} \quad (23)$$

We have also that, if $z \in D(\mathcal{L})$, then $z \in L^2_{loc}(0, \infty; D(A))$, $\frac{dz}{dt} \in L^2_{loc}(0, \infty; H)$. (By $L^2_{loc}(0, \infty; X)$ we mean the space of measurable functions $u : (0, \infty) \rightarrow X$ such that $u \in L^2(\delta, T; X)$ for all $0 < \delta < T < \infty$.) We set

$$W^{1,2}(0, \infty; H) = \{z \in L^2_{loc}(0, \infty; H); \frac{dz}{dt} \in L^2_{loc}(0, \infty; H)\}.$$

and it turns out that (see [2]) *the operator \mathcal{L} is surjective.* \square

Proof of Theorem 2. For each $f \in L^2(0, \infty; H)$, the solution $q \in L^2(0, \infty; H)$ to equation

$$\frac{dq}{dt} - \mathcal{A}^*q - (B'(y^*))^*q - (FF^*R)^*q = f, \quad t \geq 0 \quad (24)$$

is defined by

$$\langle q, \psi \rangle_{L^2(0, \infty; H)} = -\langle f, \mathcal{L}^{-1}\psi \rangle_{L^2(0, \infty; H)}, \quad \forall \psi \in L^2(0, \infty; H) \quad (25)$$

and so $\mathcal{L}^*(q) = -f$ where \mathcal{L}^* is the adjoint of \mathcal{L} .

According to this definition, the solution p to equation (19) is defined by

$$\langle p, \psi \rangle_{L^2(0, \infty; H)} = \langle A^{\frac{3}{2}}y^* - RFF^*p, \mathcal{L}^{-1}\psi \rangle_{L^2(0, \infty; H)}, \quad \forall \psi \in L^2(0, \infty; H). \quad (26)$$

Since, as seen earlier, $FF^*Rz = FF^*R\mathcal{L}^{-1}\psi \in L^2(0, \infty; H)$, (26) makes sense.

Now, coming back to problem (21), we see that for $v^*(t) = u^*(t) + F^*Ry^*(t)$ (optimal) we have

$$\int_0^\infty ((y^*(t), z(t))_{\frac{3}{2}} + (v^*(t) - F^*Ry^*(t), v(t) - F^*Rz(t))) dt = 0 \quad (27)$$

for all $v \in L^2(0, \infty; R^M)$, where z is the solution to equation

$$\mathcal{L}(z) = Fv. \quad (28)$$

(Here $(\cdot, \cdot)_{\frac{3}{2}}$ is the scalar product in $D(A^{\frac{3}{4}})$.) Then, if p is the solution to equation

$$\mathcal{L}p^* = -(A^{\frac{3}{2}}y^* - RFu^*), \quad (29)$$

we obtain by (27) that

$$\langle \mathcal{L}^*p, z \rangle_{L^2(0, \infty; H)} - \langle u^*, v \rangle_{L^2(0, \infty; H)} = 0, \quad \forall v \in L^2(0, \infty; H)$$

and we get (18) as claimed.

Theorem 3 *For $x \in \mathcal{U}_\rho$ and ρ sufficiently small the solution (y^*, u^*) to problem (16) is unique and $\varphi : W \rightarrow \mathbb{R}$ is Gâteaux differentiable on \mathcal{U}_ρ . Moreover, the semigroup $t \rightarrow y^*(t, x)$ leaves invariant the set \mathcal{U}_ρ and*

$$u^*(t) = -F^*\nabla\varphi(y^*(t)), \quad \forall t \geq 0. \quad (30)$$

Proof. The proof of uniqueness for solution (y^*, p) to system (14), (19), (20) for $x \in \mathcal{U}_\rho$ with ρ small enough follows as in [7] by standard estimates of the type used above for the solutions to system (4), (20). Moreover, by (18) we see that for all $h \in \mathcal{U}_\rho$,

$$\lim_{\lambda \downarrow 0} \frac{\varphi(x + \lambda h) - \varphi(x)}{h} = \int_0^\infty (\langle y^*(t), z(t) \rangle_{\frac{3}{2}} + (u^*(t), v(t))_{R^M}) dt$$

where (z, v) is the solution to system

$$\begin{aligned} \frac{dz}{dt} + \mathcal{A}z + B'(y^*)z &= Fv, \quad t \geq 0 \\ z(0) &= h. \end{aligned}$$

Then, we obtain that

$$- (p(0), h) = \lim_{h \downarrow 0} \frac{1}{h} (\varphi(x + \lambda h) - \varphi(x)). \quad (31)$$

Hence $-p(0) = \nabla \varphi(x)$, where p is the solution to equation (20). By the dynamic programming principle the latter implies also that (30) holds and the flow $t \rightarrow y^*(t, x)$ leaves invariant \mathcal{U}_ρ . ■

Corollary 1 *The function φ is the unique solution on \mathcal{U}_ρ to operatorial (Hamilton–Jacobi) equation*

$$(\mathcal{A}x + Bx, \nabla \varphi(x)) + \frac{1}{2} |F^* \nabla \varphi(x)|^2 = \frac{1}{2} |x|_{\frac{3}{2}}^2, \quad \forall x \in \mathcal{U}_\rho \cap D(A). \quad (32)$$

Moreover, φ is convex for a sufficiently small ρ , $\varphi(x) \geq \gamma |x|_{\frac{1}{2}}^2$, $\forall x \in \mathcal{U}_\rho$ and $(D^2 \varphi(0)h, h) = (Rh, h) \geq \gamma |h|_{\frac{1}{2}}^2$, $\forall h \in W$, where γ is a positive constant.

Proof. Equation (32) follows by (30) and the obvious relation

$$\varphi(y^*(t)) = \frac{1}{2} \int_t^\infty (|y^*(s)|_{\frac{3}{2}}^2 + |u^*(s)|_M^2) ds, \quad \forall t \geq 0. \quad (33)$$

Conversely, if φ is a solution to (32), then (33) holds and this proves uniqueness of solution φ to (32).

We note also that $D^2 \varphi(0) = R \in L(W, W') \cap L(D(A^{\frac{1}{2}}), H)$ is the solution to the algebraic Riccati equation (5) and $D^2 \varphi \in C_b(\mathcal{U}_\rho; L(D(A^{\frac{1}{2}}), H))$. In particular this implies that φ is convex in the neighborhood \mathcal{U}_ρ of the origin for ρ sufficiently small. ■

3. Stabilization by noise

Throughout in the following β_i , $i = 1, \dots$, are independent Brownian motions in a probability space $\{\Omega, \mathbb{P}, \mathcal{F}, \mathcal{F}_t\}_{t>0}$.

Consider the Navier-Stokes equation (1), i.e.,

$$\begin{aligned} X_t - \nu \Delta X + (X \cdot \nabla)X &= f_e + \nabla p \quad \text{in } D \times (0, \infty) \\ \nabla \cdot X &= 0, \quad X|_{\partial D} = 0 \\ X(0) &= x, \quad D \subset \mathbb{R}^d, \quad d \geq 2. \end{aligned} \quad (34)$$

Let X_e be a steady-state to (34), i.e. (see (2)),

$$\begin{aligned} -\nu \Delta X_e + (X_e \cdot \nabla)X_e &= f_e + \nabla p_e \quad \text{in } D \\ \nabla \cdot X_e &= 0, \quad X_e|_{\partial D} = 0. \end{aligned} \quad (35)$$

If $X \implies X - X_e$, equation (34) reduces to

$$\begin{aligned} X_t - \nu \Delta X + (X \cdot \nabla)X_e + (X_e \cdot \nabla)X + (X \cdot \nabla)X &= \nabla p \\ \nabla \cdot X &= 0, \quad X|_{\partial D} = 0 \\ X(0) &= x \end{aligned} \quad (36)$$

where X is $x - X_e$.

Or, in the space H ,

$$\begin{aligned} \dot{X}(t) + \mathcal{A}X(t) + BX(t) &= 0, \quad t \geq 0, \\ X(0) &= x. \end{aligned} \quad (37)$$

To stabilize the linearized part of (37), we associate the control stochastic problem

$$\begin{aligned} dX + \mathcal{A}X dt &= v dW_t \\ X(0) &= x \end{aligned} \quad (38)$$

where W_t is a Wiener process in a probability space $\{\Omega, \mathbb{P}, \mathcal{F}, \mathcal{F}_t\}_{t>0}$ and v a control input to be precised below.

We recall a few properties of the Stokes–Oseen operator \mathcal{A} defined above.

We shall denote by \tilde{H} the complexified space $H + iH$ with scalar product denoted $\langle \cdot, \cdot \rangle$ and norm $|\cdot|_{\tilde{H}}$. Denote again \mathcal{A} the extension of \mathcal{A} to this space. The operator \mathcal{A} has a compact resolvent $(\lambda I + \mathcal{A})^{-1}$ and $-\mathcal{A}$ generates a C_0 -analytic semigroup $e^{-\mathcal{A}t}$ in \tilde{H} . Consequently, it has a countable number of eigenvalues $\{\lambda_j\}_{j=1}^{\infty}$ with corresponding eigenfunctions φ_j each with finite algebraic multiplicity m_j .

We shall denote by N the number of eigenvalues λ_j with $\text{Re } \lambda_j \leq -\gamma$, $j = 1, \dots, N$, where γ is a fixed positive number.

Denote by \tilde{P}_N the projector on the finite dimensional space

$$\tilde{\mathcal{X}}_u = \text{lin span}\{\varphi_j\}_{j=1}^N. \quad (39)$$

We have $\tilde{\mathcal{X}}_u = \tilde{P}_N \tilde{H}$ and

$$\tilde{P}_N = \frac{1}{2\pi i} \int_{\Gamma} (\lambda I + \mathcal{A})^{-1} d\lambda,$$

where Γ is a closed smooth curve in \mathbb{C} which is the boundary of a domain containing in interior the eigenvalues $\{\lambda_j\}_{j=1}^N$.

Let $\mathcal{A}_u = P_N \mathcal{A}$, $\mathcal{A}_s = (I - P_N) \mathcal{A}$. Then $\mathcal{A}_u, \mathcal{A}_s$ leave invariant spaces $\tilde{\mathcal{X}}_u, \tilde{\mathcal{X}}_s = (I - \tilde{P}_N) \tilde{H}$ and the spectra $\sigma(\mathcal{A}_u), \sigma(\mathcal{A}_s)$ are given by (see [9])

$$\sigma(\mathcal{A}_u) = \{\lambda_j\}_{j=1}^N, \quad \sigma(\mathcal{A}_s) = \{\lambda_j\}_{j=N+1}^\infty.$$

Since $\sigma(\mathcal{A}_s) \subset \{\lambda \in \mathbb{C}; \operatorname{Re} \lambda > -\gamma\}$ and \mathcal{A}_s generates an analytic C_0 -semigroup on \tilde{H} , we have

$$|e^{-\mathcal{A}_s t} x|_{\tilde{H}} \leq C e^{-\gamma t} |x|_{\tilde{H}}, \quad \forall x \in \tilde{H}, t \geq 0. \quad (40)$$

We set

$$\psi_j^1 = \operatorname{Re} \varphi_j, \quad \psi_j^2 = \operatorname{Im} \varphi_j, \quad j = 1, \dots, N, \quad (41)$$

and

$$\begin{aligned} \mathcal{X}_u &= \operatorname{lin} \operatorname{span} \{ \{\psi_j^1\} \cup \{\psi_j^2\} \}_{j=1}^N \\ \mathcal{X}_s &= \operatorname{lin} \operatorname{span} \{ \{\psi_j^1\} \cup \{\psi_j^2\} \}_{j=N+1}^\infty. \end{aligned}$$

Clearly, \mathcal{A} leaves invariant the real spaces \mathcal{X}_u and \mathcal{X}_s . More precisely, we have

$$\tilde{\mathcal{X}}_u = \mathcal{X}_u + i\mathcal{X}_u, \quad \tilde{\mathcal{X}}_s = \mathcal{X}_s + i\mathcal{X}_s$$

and therefore, $H = \mathcal{X}_u \oplus \mathcal{X}_s$, the direct sum of \mathcal{X}_u and \mathcal{X}_s (see, e.g., [9]). Since the system $\{\lambda_j\}_{j=1}^N$ is of the form

$$\{\xi_j \pm i\eta_j\}_{j=1}^{2M}, \quad \{\delta_j\}_{j=1}^{M_0}, \quad 2M + M_0 = N, \quad \xi_j, \eta_j, \delta_j \in R$$

it follows that

$$\mathcal{X}_u = \operatorname{lin} \operatorname{span} \{ \psi_j \}_{j=1}^N, \quad (42)$$

where

$$\begin{aligned} \psi_j &= \psi_j^1, \quad 1 \leq j \leq M, \quad \psi_j = \psi_j^2, \quad M < j \leq 2M, \\ \psi_j &= \varphi_j \text{ (real eigenfunctions)}, \quad 2M < j \leq N. \end{aligned} \quad (43)$$

The infinite dimensional space \mathcal{X}_s is similarly generated and estimate (40) remains valid in the $|\cdot|_N$ -norm for \mathcal{A}_s defined on $\mathcal{X}_s \subset H$. We shall denote by P_N the projector corresponding to the decomposition $H = \mathcal{X}_u \oplus \mathcal{X}_s$, i.e.,

$$\mathcal{X}_u = P_N H, \quad \mathcal{X}_s = (I - P_N) H.$$

Consider the orthonormal system $\{\phi_j\}_{j=1}^N$ obtained from $\{\psi_j\}_{j=1}^N$ by the Schmidt orthogonalization procedure.

Consider the following stochastic perturbation of the linearized system (5) (see (6))

$$\begin{aligned} dX + \mathcal{A}X dt &= \eta \sum_{i=1}^N (X(t), \phi_i) P(m\phi_i) d\beta_i(t) \\ X(0) &= x, \end{aligned} \quad (44)$$

where $\eta \in R$ and $m = \chi_{\mathcal{O}_0}$ is the characteristic of the open subset $\mathcal{O}_0 \subset D$.

Equation (44) can be seen as a closed loop system associated to the controlled equation

$$\begin{aligned} dX + \mathcal{A}X dt &= \eta \sum_{i=1}^N v_i(t) P(m\phi_i) d\beta_i(t), \quad t \geq 0, \\ X(0) &= x, \end{aligned} \tag{45}$$

with the feedback controller $v_i = (X(t), \phi_i)$, $i = 1, \dots, N$.

Theorem 4 *Let $X = X(t, x)$ be the solution to (45). Then, for $|\eta|$ sufficiently large, we have*

$$\mathbb{P} \left\{ \lim_{t \rightarrow \infty} e^{\gamma t} X(t, x) = 0 \right\} = 1, \quad \forall x \in H, \tag{46}$$

The closed loop system (44) can be equivalently written as

$$\begin{aligned} dX(t) - \nu \Delta X(t) dt + (X(t) \cdot \nabla) X_e dt + (X_e \cdot \nabla) X(t) dt \\ = \eta m \sum_{i=1}^N (X(t), \phi_i) \phi_i d\beta_i(t) + \nabla p(t) dt \text{ in } (0, \infty) \times D \\ \nabla \cdot X(t) = 0, \quad X(t)|_{\partial D} = 0 \\ X(0) = x \text{ in } D. \end{aligned} \tag{47}$$

In particular, it follows by Theorem 4 that the feedback controller $u_i = \eta m (X - X_e, \phi_i) \phi_i$, $i = 1, \dots, N$, stabilizes exponentially the stationary solution X_e , i.e., we have

Corollary 2.1. *The solution X to closed loop system*

$$\begin{aligned} dX(t) - \nu \Delta X(t) dt &= \eta m \sum_{i=1}^N (X(t) - X_e, \phi_i) \phi_i d\beta_i(t) + \nabla p dt, \quad t \geq 0 \\ \nabla \cdot X &= 0, \quad X|_{\partial D} = 0 \\ X(0) &= x \end{aligned} \tag{48}$$

satisfies

$$\mathbb{P} \left[\lim_{t \rightarrow \infty} (X(t) - X_e) e^{\gamma t} = 0 \right] = 1, \quad \forall x \in H. \tag{49}$$

The proof of Theorem 4 follows from same sharp arguments involving the martingale theory and is given in [3].

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